## Complex Gaussian multiplicative chaos

#### Vincent Vargas (with H. Lacoin and R. Rhodes)

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3 Motivations from Liouville Quantum Gravity

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We want to give a precise meaning to distributions  $M^{\gamma,\beta}$  defined formally by:

$$M^{\gamma,\beta}(A) = \int_A e^{\gamma X(x) + i\beta Y(x)} dx, \ A \subset D$$

where:

- $\gamma, \beta \geq 0$
- $D \subset \mathbb{C}$  a bounded domain
- *X*, *Y* two centered **independent** GFF with covariance given by:

$$E[X(x)X(y)] = G_D(x,y) \underset{|x-y| 
ightarrow 0}{\sim} \ln rac{1}{|x-y|}$$

where  $G_D$  is the Green kernel:

 $-\Delta_{y}G_{D}(x,y)=2\pi\delta_{x}$ 

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We consider a family of centered Gaussian processes  $(X_{\varepsilon}(x))_{x \in D}$  $(\varepsilon \leq 1)$ :

• Covariance: 
$$E[X_{\varepsilon}(x)X_{\varepsilon}(y)] \sim \ln \frac{1}{|x-y|+\varepsilon} \underset{\varepsilon \to 0}{\to} G_D(x,y)$$

Variance: E[X<sub>ε</sub>(x)<sup>2</sup>] = ln <sup>1</sup>/<sub>ε</sub> + ln C(x, D) + o(1) where C(x, D) conformal radius.

•  $\varepsilon \mapsto X_{\varepsilon}$  independent increments

Same for  $(Y_{\varepsilon}(x))_{x \in D}$  ( $\varepsilon \leq 1$ ) independent from X.

We define:

$$M^{\gamma,\beta}_{\varepsilon}(A) = \int_{A} e^{\gamma X_{\varepsilon}(x) + i\beta Y_{\varepsilon}(x)} dx, \ A \subset D$$

Observe that:

$$M^{\gamma,0}_{arepsilon}(A) = \int_{A} e^{\gamma X_{arepsilon}(x)} dx, \ A \subset D$$

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One can also work with other "smooth" cut-offs

- 1985: Kahane, H<sup>1</sup>-basis decomposition
- 2006-2008: Robert, V., general convolutions
- 2008: Duplantier, Sheffield , circle averages

In fact, one can work with any log-correlated field in any dimension (Kahane, 1985, Robert, V., 2006, 2008): see our review with Rhodes.

#### Theorem (Kahane, 1985)

There exists a random measure  $M^{\gamma,0}$  such that following limit exists almost surely in the space of Radon measures:

$$\varepsilon^{\frac{\gamma^2}{2}}M^{\gamma,0}_{\varepsilon}(dx) \xrightarrow[\varepsilon o 0]{} M^{\gamma,0}(dx).$$

 $M^{\gamma,0}$  is called Gaussian multiplicative chaos associated to the Green kernel.

Remark

When J.P. Kahane meets Paul Levy ...

## Gaussian multiplicative chaos: $\beta = 0$

#### Theorem (Kahane, 1985)

The measure  $M^{\gamma,0}$  is different from 0 if and only if  $\gamma < 2$ .

#### Theorem (Kahane, 1985)

For  $\gamma < 2$ , the measure  $M^{\gamma,0}$  "lives" almost surely on a set of Hausdorff dimension  $2 - \frac{\gamma^2}{2}$  (the set of  $\gamma$ -thick points).

See also Hu, Miller, Peres (2010).

## Density of Gaussian multiplicative chaos with respect to $\gamma$

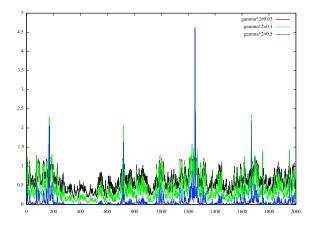
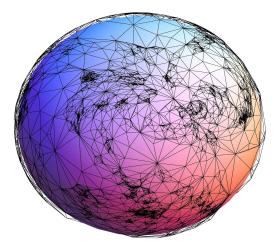


Figure: Density of Gaussian multiplicative chaos

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# Uniformisation of a uniform triangulation: courtesy of N. Curien



It is conjectured to be the limit of random planar maps weighted by a critical statistical physics system (CFT with central charge  $c \leq 1$ ) and conformally mapped to a domain D:

- Ambjorn-Durhuus-Jonsson (2005): Quantum geometry: A Statistical Field Theory Approach.
- Sheffield (2010): Conformal weldings of random surfaces: SLE and the quantum gravity zipper. Precise math conjectures.

- Curien (2013): A glimpse of the conformal structure of random planar maps (c = 0). First step in a mathematical proof.
- Miller, Sheffield (2014): Quantum Loewner evolution.

#### Theorem (Duplantier, Rhodes, Sheffield, V., 2012)

There exists a random measure M such that following limit exists almost surely in the space of Radon measures:

$$arepsilon^2(2\lnrac{1}{arepsilon}-X_arepsilon(x))M^{2,0}_arepsilon(dx) \stackrel{
ightarrow}{
ightarrow} M^{'}(dx).$$

M' is called critical Gaussian multiplicative chaos associated to the Green kernel.

#### Theorem (Duplantier, Rhodes, Sheffield, V., 2012)

The following limit exists almost surely (along suitable subsequences) in the space of Radon measures:

$$\sqrt{\ln \frac{1}{\varepsilon}} \varepsilon^2 M_{\varepsilon}^{2,0}(dx) \underset{\varepsilon \to 0}{\to} \sqrt{\frac{2}{\pi}} M'(dx).$$

Theorem (Barral, Kupiainen, Nikula, Saksman, Webb, 2013)

The measure M' lives on a set of Hausdorff dimension 0.

## Complex Gaussian multiplicative chaos: Phase diagram

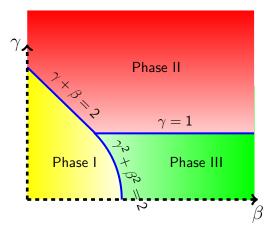


Figure: Phase diagram

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Previous work on the complex case:

• Computation of the Free Energy of complex multiplicative cascades: Derrida, Evans, Speer, 1993. In our context:

$$\lim_{\varepsilon \to 0} \frac{\ln |M_{\varepsilon}^{\gamma,\beta}([0,1]^2)|}{\ln \frac{1}{\varepsilon}}$$

• Complex multiplicative cascades: series of works in dimension 1 by Barral, Jin, Mandelbrot, 2010. Essentially investigated phase I. Partial results in phase III.

## Convergence in phase I and it's frontier I/II (excluding extremal points)

#### Theorem (Lacoin, Rhodes, V., 2013)

On phase I and it's frontier I/II (excluding the extremal points), the D'(D)-valued distribution:

$$M_{\varepsilon}^{\gamma,\beta}: \varphi \to \varepsilon^{\frac{\gamma^2}{2} - \frac{\beta^2}{2}} \int_D \varphi(x) M_{\varepsilon}^{\gamma,\beta}(dx)$$

converges almost surely in the space  $\mathcal{D}'_2(D)$  of distributions of order 2 towards a non trivial limit  $M^{\gamma,\beta}$ .

The probability measure is  $e^{-S(Y)}dY$  where S(Y) is the action:

$$S(Y) = \frac{1}{4\pi} \int_D |\nabla Y(x)|^2 dx + \mu \int_D \cos(\beta Y(x)) dx$$

Different regimes ( $\varepsilon \rightarrow 0$ ):

- $\beta^2 < 2$ : non trivial convergence of  $\mathbb{E}[e^{-\mu\varepsilon^{-\beta^2/2}\int_D \cos(\beta Y_{\varepsilon}(x))dx}]$ •  $2 \leq \beta^2 < \beta_c^2$ :  $\varepsilon^{-\beta^2/2}\int_D \cos(\beta Y_{\varepsilon}(x))dx \approx \sigma_{\varepsilon}N + O(1)$  where  $\sigma_{\varepsilon} \to \infty$  and N Gaussian variable.
- $\beta^2 > \beta_c^2$ : more and more cumulants to take out...

## Convergence in the inner phase III and it's frontier I/III

#### Theorem (Lacoin, Rhodes, V., 2013)

• When  $\gamma \in [0,1[$  and  $\beta^2 + \gamma^2 > 2$ , we have

$$\left(\varepsilon^{\gamma^2-1}M_{\varepsilon}^{\gamma,\beta}(A)\right)_{A\subset D} \Rightarrow (W_{\sigma^2M^{2\gamma,0}}(A))_{A\subset D}.$$
 (1)

where  $\sigma^2 := \sigma^2(\beta^2 + \gamma^2) > 0$  and W is a complex Gaussian measure on D with intensity  $\sigma^2 M^{2\gamma,0}$ .

• When  $\gamma \in [0,1[$  and  $\beta^2 + \gamma^2 = 2$ , we have

$$\left(\varepsilon^{\gamma^2-1}|\log\varepsilon|^{-1/2}M_{\varepsilon}^{\gamma,\beta}(A)\right)_{A\subset D} \Rightarrow (W_{\sigma^2M^{2\gamma,0}}(A))_{A\subset D}.$$
 (2)

where  $\sigma^2 > 0$  and W is a complex Gaussian measure on D with intensity  $\sigma^2 M^{2\gamma,0}$ .

#### Theorem (Lacoin, Rhodes, V., 2013)

When  $\gamma = 1$  and  $\beta^2 + \gamma^2 > 2$ , we have

$$\left(|\ln \varepsilon|^{1/4} M_{\varepsilon}^{\gamma,\beta}(A)\right)_{A\subset D} \Rightarrow (W_{\sigma^2 M'}(A))_{A\subset D}.$$

with  $\sigma^2 := \sigma^2(\beta) > 0$  and  $W_{\sigma^2 M'}(\cdot)$  is a complex Gaussian random measure with intensity  $\sigma^2 M'$ .

## Conformal Field theory c = 1 coupled to Liouville Quantum Gravity

Recall that, on phase I and it's frontier I/II (excluding the extremal points), the  $\mathcal{D}'(D)$ -valued distribution:

$$M_{\varepsilon}^{\gamma,\beta}:\varphi\to\varepsilon^{\frac{\gamma^{2}}{2}-\frac{\beta^{2}}{2}}\int_{D}\varphi(x)M_{\varepsilon}^{\gamma,\beta}(dx)$$

converges almost surely in the space  $\mathcal{D}'_2(D)$  of distributions of order 2 towards a non trivial limit  $M_{X,Y}^{\gamma,\beta}$ .

In fact, we must denote

$$M_{X,Y}^{\gamma,\beta}(dx) = e^{\gamma X(x) + i\beta Y(x) - \frac{\gamma^2}{2} \mathbb{E}[X(x)^2] + \frac{\beta^2}{2} \mathbb{E}[Y(x)^2]} C(x,D)^{\frac{\gamma^2}{2} - \frac{\beta^2}{2}} dx,$$

where C(x, D) is the conformal radius. This is because we do not renormalize by the mean!

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Polyakov action on a domain D

$$S(X,Y) = \frac{1}{4\pi} \int_{D} |\nabla Y(x)|^2 dx + \frac{1}{4\pi} \int_{D} |\nabla X(x)|^2 + QR(x)X(x)dx,$$

R is the curvature and Q = 2

• Equivalence class of random surfaces:

$$(X, Y) \rightarrow (X \circ \psi + 2 \ln |\psi'|, Y \circ \psi),$$

where  $\psi : \tilde{D} \to D$  is a conformal map. See Ginsparg, Moore (1993), Lectures on 2D gravity and 2D string theory.

Under the above equivalence class  $(\psi: \tilde{D} \rightarrow D)$ 

$$M_{X\circ\psi+2\ln|\psi'|,Y\circ\psi}^{\gamma,\beta}(\varphi)=|\psi'\circ\psi^{-1}|^{2\gamma-\frac{\gamma^2}{2}+\frac{\beta^2}{2}-2}M_{X,Y}^{\gamma,\beta}(\varphi\circ\psi^{-1}),$$

for every function  $\varphi \in C^2_c(\tilde{D})$ 

Tachyon Fields are conformally invariant. One must solve

$$2\gamma - \frac{\gamma^2}{2} + \frac{\beta^2}{2} - 2 = 0 \leftrightarrow \gamma \pm \beta = 2, \ \gamma \in ]1, 2[.$$

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At the special point ( $\gamma = 2, \beta = 0$ ), we recover the special tachyon field, i.e. the background measure

$$M^{\gamma,\beta}_{X,Y}(A) = M'(A)$$

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where M' is critical Gaussian multiplicative chaos.

The probability measure is  $e^{-S(Y)}dY$  where S(Y) is the action:

$$S(Y) = \frac{1}{4\pi} \int_D |\nabla Y(x)|^2 dx + \mu \int_D \cos(\beta Y(x)) dx$$

Representation of the density of charge  $\rho$  of the Coulomb gas:

$$<
ho(x)
ho(y)>=\mathbb{E}[\sin(eta Y(x))\sin(eta Y(y))e^{-\mu\int_D\cos(eta Y(z))dz}]$$

## Sine-Gordon model coupled to gravity?

The probability measure is  $e^{-S(X,Y)}dXdY$  where S(X,Y) is the action:

$$S(X,Y) = \frac{1}{4\pi} \int_{D} |\nabla Y(x)|^2 dx + \frac{1}{4\pi} \int_{D} |\nabla X(x)|^2 dx + \mu_1 \int_{D} \cos(\beta Y(x)) e^{\gamma X(x)} dx + \mu_2 \int_{D} e^{2X(x)} dx$$

where  $\gamma + \beta = 2$  (see G. Moore, Gravitational Phase transitions and the Sine-Gordon model).

The problem is linked to defining the Coulomb gas on a random lattice.