

# High-frequency Helmholtz problems: Lecture 1

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Webpage for the lectures:

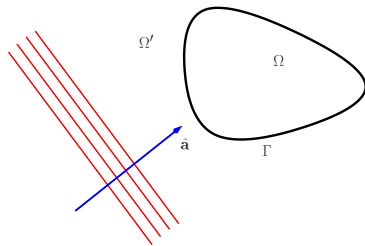
<http://people.bath.ac.uk/masigg/CUHK/webpage/>

S.N. Chandler-Wilde, I.G.G., S.Langdon, E.A. Spence, Numerical-asymptotic  
boundary integral methods in high-frequency scattering  
*Acta Numerica* 2012, pp 85-305

# High freq. problem for the Helmholtz equation

Given an object  $\Omega \subset \mathbb{R}^d$ , with boundary  $\Gamma$  and exterior  $\Omega'$ ,

**Incident plane wave, e.g. :**  $u_I(x) = \exp(i\mathbf{k}\mathbf{x} \cdot \hat{\mathbf{a}})$



**Total wave**  $u = u_I + u_S$ , where **Scattered wave**  $u_S$  satisfies:

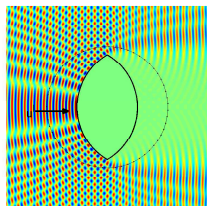
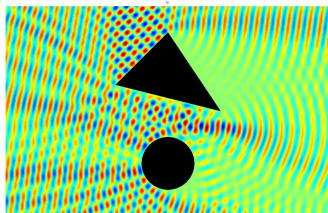
$$\Delta u_S + \mathbf{k}^2 u_S = 0 \quad \text{in } \Omega'$$

plus **boundary condition** (Mostly  $u_I + u_S = 0$  on  $\Gamma$ ) and

**radiation condition:**  $\frac{\partial u^S}{\partial r} - i\mathbf{k}u^S = o(r^{-(d-1)/2})$  as  $r \rightarrow \infty$

- Oscillatory solutions
- Complexity:  $\mathcal{O}(k^d)$  FEM ,  $\mathcal{O}(k^{d-1})$  BEM
- “Pollution effect”?
- BEM suitable for homogeneous problems, i.e. (piecewise) constant wavenumbers
- **Lecture 1: wavenumber dependent NA for BEM**  
[§6, ChGrLaSp] [GrLoMeSp]

# Numerical-asymptotic methods



**Lecture 2:** By building in asymptotic information about solution we can reduce (or remove) the wavenumber dependence.

**Interesting mathematics** [ChLa], [DoGrSm], [ChGrLaSp]

**BUT**

**Methods are strongly geometry dependent**

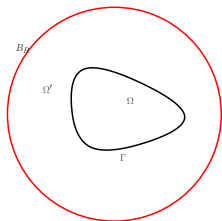
# Truncated problems

$$\Delta u_S + k^2 u_S = 0 \quad \text{in } \Omega' \cap B_R$$

$$u_S = -u_I \quad \text{on } \Gamma$$

$$\frac{\partial u_S}{\partial n} - iku_S = 0 \quad \text{on } B_R$$

for large  $R$



**Model “cavity” problem:**  $u_S \rightarrow u$

$$\Delta u + k^2 u = f \quad \text{in bounded domain } \Omega$$

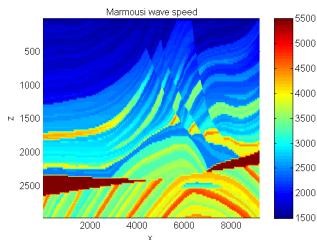
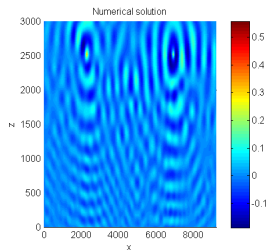
$$\frac{\partial u}{\partial n} - iku = g \quad \text{on } \Gamma := \partial\Omega$$

# Heterogeneity

Seismic inversion problem:

$$-\Delta u - \left( \frac{\omega L}{c(x)} \right)^2 u = f, \quad \omega = \text{frequency}$$

solve for  $u$  with approximate  $c$ .



Third talk: Conventional discretisation and fast solvers

# First problem

When (i.e. for what values of  $h$ ) is the error in the  $h$ -version boundary element method (BEM) bounded independently of  $k$ ?

First: Short description of BEM. [ChGrLaSp]

# Fundamental solution for the Helmholtz equation

## Blackboard

$$-(\Delta u + k^2 u) = 0$$

$$G_k(x, y) = \begin{cases} \frac{i}{4} H_0^{(1)}(k|x-y|) & \text{2D} \\ \frac{\exp(ik|x-y|)}{4\pi|x-y|} & \text{3D} \end{cases}$$

Phase:  $k|x-y| \implies$  **Oscillatory integral**

**single layer potential** :  $(\mathcal{S}_k \phi)(x) = \int_{\Gamma} G_k(x, y) \phi(y) dS(y),$

**double layer**:  $(\mathcal{D}_k \phi)(x) = \int_{\Gamma} [\partial_{n(y)} G_k(x, y)] \phi(y) dS(y),$

**adjoint double layer**:  $\mathcal{D}'_k$  (switch roles of  $x$  and  $y$ ).



# Combined potential boundary integral formulations

Exterior scattering problem with incident field  $u_I$ :

Green's identity for  $u_S$  in  $\Omega'$ :

$$\mathcal{S}_k(\partial_n u_S) - \mathcal{D}_k(u_S) = (-u_S) \quad \text{in } \Omega' \quad (1)$$

# Combined potential boundary integral formulations

Exterior scattering problem with incident field  $u_I$ :

Green's identity for  $u_I$  in  $\Omega$ :

$$\mathcal{S}_k(\partial_n u_S + \partial_n u_I) - \mathcal{D}_k(u_S + u_I) = (-u_S + 0) \quad \text{in } \Omega' \quad (1)$$

# Combined potential boundary integral formulations

Exterior scattering problem with incident field  $u_I$ :

Green's identity for  $u_I$  in  $\Omega$ :

$$\mathcal{S}_k(\underbrace{\partial_n u_S + \partial_n u_I}_{\partial_n u}) - \mathcal{D}_k(\underbrace{u_S + u_I}_{=0}) = \underbrace{(-u_S + 0)}_{u_I} \quad \text{in } \Omega' \quad (1)$$

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Limit to boundary  $\Gamma$ : Equation for unknown  $v := \partial_n u$   
but with spurious frequencies.

$k$

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Limit to boundary  $\Gamma$ : Equation for unknown  $v := \partial_n u$   
but with spurious frequencies.

Take normal derivative in (1) and combine with  $-ik \times$  (1):

**“direct” combined potential formulation** **Blackboard**

$$\mathcal{R}'_k v := \left( \frac{1}{2}I + \mathcal{D}'_k \right) v - ik \mathcal{S}_k v = \partial_n u_I - ik u_I, \quad \text{or } k \rightarrow \eta$$

# Combined potential boundary integral formulations

Exterior scattering problem with incident field  $u_I$ :

Green's identity for  $u_I$  in  $\Omega$ :

$$\mathcal{S}_k(\underbrace{\partial_n u_S + \partial_n u_I}_{\partial_n u}) - \mathcal{D}_k(\underbrace{u_S + u_I}_{=0}) = \underbrace{(-u_S + 0)}_{u_I} \quad \text{in } \Omega' \quad (1)$$

Limit to boundary  $\Gamma$ : Equation for unknown  $v := \partial_n u$  but with spurious frequencies.

Take normal derivative in (1) and combine with  $-ik \times (1)$ :

**“direct” combined potential formulation**

$$\mathcal{R}'_k v := \left( \frac{1}{2}I + \mathcal{D}'_k \right) v - ik \mathcal{S}_k v = \partial_n u_I - ik u_I ,$$

Alternative **“indirect”** method:

$$\mathcal{R}_k \phi := \left( \frac{1}{2}I + \mathcal{D}_k \right) \phi - ik \mathcal{S}_k \phi = -u_I ,$$

“Fredholm integral equations of the Second kind”

$$\begin{aligned}\mathcal{R}'_k v &= (\lambda I + \mathcal{L}'_k)v = f_k \\ \mathcal{R}_k \phi &= (\lambda I + \mathcal{L}_k)\phi = g_k \quad (\lambda = 1/2)\end{aligned}$$

Galerkin method in approximating space  $\mathcal{V}_N$  (or  $\mathcal{V}_h$ ).

e.g. piecewise polynomials of fixed degree  $p$ . **Blackboard**

Solution  $v_N$  or  $\phi_N$ , e.g.

$$(\lambda I + \mathcal{P}_N \mathcal{L}'_k)v_N = \mathcal{P}_N f_k$$

“Fredholm integral equations of the second kind”

$$\begin{aligned}\mathcal{R}'_k v &= (\lambda I + \mathcal{L}'_k)v = f_k \\ \mathcal{R}_k \phi &= (\lambda I + \mathcal{L}_k)\phi = g_k \quad (\lambda = 1/2)\end{aligned}$$

Galerkin method in approximating space  $\mathcal{V}_N$  (or  $\mathcal{V}_h$ ).

e.g. piecewise polynomials of fixed degree  $p$ . **Blackboard**

Solution  $v_N$  or  $\phi_N$ , e.g.

$$(\lambda I + \mathcal{P}_N \mathcal{L}'_k)v_N = \mathcal{P}_N f_k$$

$$v - v_N = \lambda \underbrace{(\lambda I - \mathcal{P}_N \mathcal{L}'_k)^{-1}}_{\text{stability}} \underbrace{(v - \mathcal{P}_N v)}_{\text{best approx}}$$



# Question 1 (best approximation error)

When are

$$\frac{\inf_{w_N \in \mathcal{V}_N} \|v - w_N\|_{L^2(\Gamma)}}{\|v\|_{L^2(\Gamma)}}$$

and

$$\frac{\inf_{w_N \in \mathcal{V}_N} \|\phi - w_N\|_{L^2(\Gamma)}}{\|\phi\|_{L^2(\Gamma)}}$$

bounded independently of  $k$ ?

## Question 2 (quasioptimality)

When are

$$\frac{\|v - v_N\|_{L^2(\Gamma)}}{\inf_{w_N \in \mathcal{V}_N} \|v - w_N\|_{L^2(\Gamma)}}$$

and

$$\frac{\|\phi - \phi_N\|_{L^2(\Gamma)}}{\inf_{w_N \in \mathcal{V}_N} \|\phi - w_N\|_{L^2(\Gamma)}}$$

bounded independently of  $k$ ?

**“Pollution effect”?**

## If both hold...(bound on relative errors)

$$\frac{\|v - v_N\|_{L^2(\Gamma)}}{\|v\|_{L^2(\Gamma)}}$$

and

$$\frac{\|\phi - \phi_N\|_{L^2(\Gamma)}}{\|\phi\|_{L^2(\Gamma)}}$$

bounded independently of  $k$ .

When is

$$\frac{\inf_{w_N \in \mathcal{V}_N} \|v - w_N\|_{L^2(\Gamma)}}{\|v\|_{L^2(\Gamma)}}$$

bounded independently of  $k$ ?

**Theorem** If  $\Omega$  is  $C^\infty$  and convex then for  $h$ -BEM,

$$\inf_{w_h \in \mathcal{V}_h} \|v - w_h\|_{L^2(\Gamma)} \lesssim (hk)^p \|v\|_{L^2(\Gamma)}$$

so  $hk \lesssim 1$  is sufficient for Question 1.

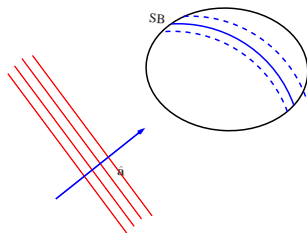
# Proof uses... Melrose and Taylor formula (1985)

$$v(\mathbf{x}) := \partial u / \partial n(\mathbf{x}) = kV(\mathbf{x}, k) \exp(ik\mathbf{x} \cdot \hat{\mathbf{a}}), \quad x \in \Gamma,$$

**Theorem** Dominguez, IGG, Smyshlyaev, 2007

$$|D^n V(x, k)| \leq \begin{cases} C_n, & n = 0, 1, \\ C_n k^{(n-1)/3} (1 + k^{1/3} \text{dist}(x, SB))^{-(n+2)} & n \geq 2, \end{cases}$$

where  $SB = \{\mathbf{x} \in \Gamma : \mathbf{n}(\mathbf{x}) \cdot \hat{\mathbf{a}} = 0\}$  shadow boundary.



**Proves, e.g.**  $\|v\|_{H^1(\Gamma)} \lesssim k \|v\|_{L^2(\Gamma)}$

# Answers: Question 1 (“direct” version $v = \partial_n u$ )

When is

$$\frac{\inf_{w_N \in \mathcal{V}_N} \|v - w_N\|_{L^2(\Gamma)}}{\|v\|_{L^2(\Gamma)}}$$

bounded independently of  $k$ ?

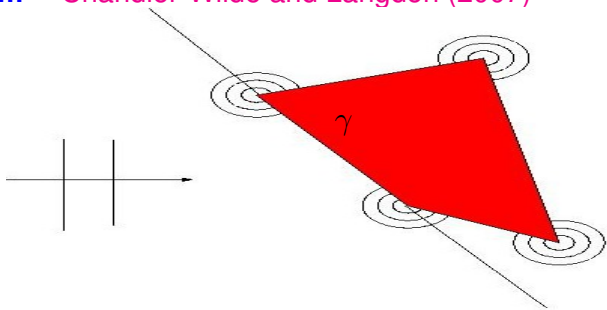
**Theorem** If  $\Omega$  is a convex polygon then there is a mesh with  $\mathcal{O}(N)$  points so that ,

$$\inf_{w_h \in \mathcal{V}_h} \|v - w_h\|_{L^2(\Gamma)} \lesssim \frac{k}{N} \|v\|_{L^2(\Gamma)}$$

so  $k/N \lesssim 1$  is sufficient for Question 1.

(Requires  $\sup_{\mathbf{x} \in \Omega'} |u(\mathbf{x})| < \infty$ .)

## Theorem Chandler-Wilde and Langdon (2007)



$$v(s) = \frac{\partial u}{\partial n}(s) = 2 \frac{\partial u^I}{\partial n}(s) + e^{iks} v_+(s) + e^{-iks} v_-(s)$$

where  $s$  is distance along  $\gamma$ , and

$$k^{-n} |v_+^{(n)}(s)| \leq \begin{cases} C_n (ks)^{-1/2-n}, & ks \geq 1, \\ C_n (ks)^{-\alpha-n}, & 0 < ks \leq 1, \end{cases}$$

where  $\alpha < 1/2$  depends on the corner angle.

# Answers: Question 1: Indirect method

$$\lambda\phi = \mathcal{L}_k\phi = ik\mathcal{S}_k\phi + \mathcal{D}_k\phi$$

To estimate the derivatives of  $\phi$ :

$$\|\mathcal{S}_k\|_{H^1 \leftarrow L^2} \lesssim k^{(d-1)/2} \quad (\Gamma \text{ Lipschitz})$$

$$\|\mathcal{D}_k\|_{H^1 \leftarrow L^2} \lesssim k^{(d+1)/2} \quad (\Gamma \text{ smooth enough})$$

These imply  $\|\phi\|_{H^1(\Gamma)} \lesssim k^{(d+1)/2} \|\phi\|_{L^2(\Gamma)}$

And so  $hk^{(d+1)/2} \lesssim 1$  is sufficient for Question 1.



# Answers: Question 2 (classical approach)

$$\begin{aligned}\mathcal{R}'_k v &:= (\lambda I + \mathcal{L}'_k)v &&= f_k \quad \text{compact perturbation} \\ &(\lambda I + \mathcal{P}_h \mathcal{L}'_k)v_h &&= \mathcal{P}_h f_k \quad \text{Galerkin method}\end{aligned}$$

**Lemma** [Atkinson, Anselone, 1960's ....]

$$\begin{aligned}\text{If} \quad & \| (I - \mathcal{P}_h) \mathcal{L}'_k \| \| (\lambda I + \mathcal{L}'_k)^{-1} \| \ll 1, \\ \text{then} \quad & \| v - v_h \| \lesssim \| (\lambda I + \mathcal{L}'_k)^{-1} \| \inf_{w_h \in \mathcal{V}_h} \| v - w_h \|\end{aligned}$$

**Application:**

$$\| (I - \mathcal{P}_h) \mathcal{L}'_k \| \lesssim h \| \mathcal{L}'_k \|_{L^2 \rightarrow H^1} \lesssim h k^{(d+1)/2}$$

**and in addition:**

$$\| (\lambda I + \mathcal{L}'_k)^{-1} \| \lesssim 1 \quad [\text{Chandler-Wilde \& Monk, 2008}]$$

Lipschitz star-shaped

**Theorem** Hence quasioptimality if  $h k^{(d+1)/2} \leq C$

## We used in this talk

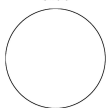
- $k$ – explicit bounds on norms of  $\mathcal{L}_k, \mathcal{L}'_k$   
(where  $\mathcal{R}'_k = \frac{1}{2}I + \mathcal{L}'_k$ ), etc.  
needed smooth enough domains
- $k$ – explicit bounds on inverses  $(\mathcal{R}_k)^{-1}, (\mathcal{R}'_k)^{-1}$   
needed Lipschitz star-shaped

# The Subtlety of Behaviour of $\|\mathcal{L}_k\|$ and $\|\mathcal{R}_k^{-1}\|$ Equivalently $\|\mathcal{L}'_k\|$ and $\|(\mathcal{R}'_k)^{-1}\|$

$$\|\mathcal{L}_k\|, \|\mathcal{R}_k^{-1}\|$$

$$\sim k^{1/3}, \sim 1$$

Circle



Ellipse

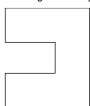


$$\sim k^{1/2}, \sim 1$$

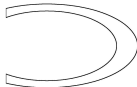
Square



Rectangular cavity



Elliptic Cavity



$$\sim k_m^{1/2}, \sim k_m^{7/5}$$

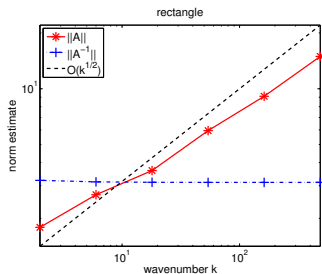
$$\sim k_m^{1/2}, \sim e^{\gamma k_m}$$

Proofs are in [ChGrLaLi], [BeChGrLaLi]

# Numerical Experiments: domain $[0, 0.5] \times [0, 5]$

$$\sqrt{1 + \gamma_p^2} \approx \frac{\|v - v_N\|_{L^2(\Gamma)}}{\inf_{w_N \in \mathcal{V}_N} \|v - w_N\|_{L^2(\Gamma)}}$$

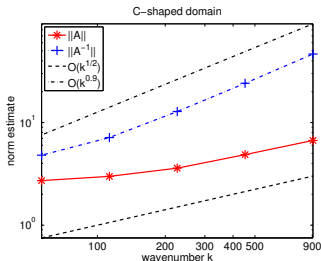
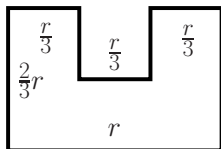
degree  $p = 0, 1$



$N$	$k$	$\gamma_0$	$\gamma_1$
22	2	0.368234	0.136623
66	6	0.334368	0.121106
198	18	0.337487	0.120028
594	54	0.335113	0.120023
1782	162	0.333687	0.12
5346	486	0.333559	0.119998

$$hk \sim 1$$

# Numerical Experiments: trapping domain



$m$	$k$	$N$	$\gamma_0$	$\gamma_1$
3	56.5	120	0.480033	0.174585
6	113.1	240	0.487655	0.174454
12	226.2	480	0.51861	0.174301
24	452.4	960	0.527743	0.174264
48	904.8	1920	0.549879	0.174278

**Open question: Prove  $hk \lesssim 1$  sufficient for quasioptimality**

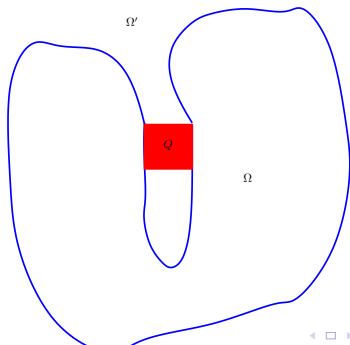
# “Trapping domains”

Can things get bad in the non-star-shaped case?

## Theorem

If the exterior domain  $\Omega'$  contains a square  $Q$  of side length  $a$  and the boundary  $\Gamma$  coincides with two parallel sides of  $Q$ , then if  $2ak = m\pi$  for any positive integer  $m$ ,

$$\|\mathcal{R}_k^{-1}\| \gtrsim (ak)^{9/10}.$$



# Non-star shaped case - “Quasimodes”

(family of) sources  $f$  and solutions  $v$  of

$$\Delta v + k^2 v = f \quad \text{in } \Omega' \quad \text{with } v = 0 \quad \text{on } \Gamma$$

+ Sommerfeld radiation condition, where

$$\|v\|_{L^2(\Omega')} \geq M_k \|f\|_{L^2(\Omega')}, \quad \text{with } M_k \text{ “large”}$$

. This would contradict the stability bound

$$\left\{ |v|_{H^1(\Omega')}^2 + k^2 \|v\|_{L^2(\Omega')}^2 \right\}^{1/2} \lesssim \|f\|_{L^2(\Omega')}$$

which holds in star-shaped case (see Lectures 3 and 4).

**Application to BIE case: Blackboard**

$$\mathcal{R}'_k (\partial_n v) = (\partial_n v^N - ikv^N)$$

where  $v^N$  is the Newtonian potential generated by  $f$

implies growth of  $\|(\mathcal{R}'_k)^{-1}\|$

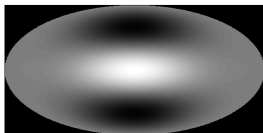
More generally ...

$$\|(\mathcal{R}'_k)^{-1}\| \gtrsim k^{-(d-2)} M_k - \mathcal{O}(k^{(d-1)/2})$$

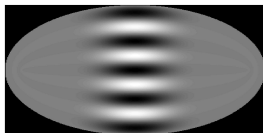
With elliptic cavity  $M_k$  can increase exponentially.

[BeChGrLaLi]

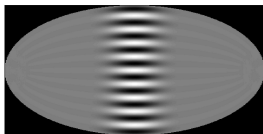
$$k_{1,0} = 9.9771201566136298$$



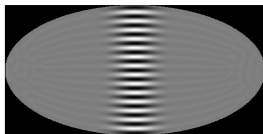
$$k_{4,0} = 28.807002784875433$$



$$k_{9,0} = 60.218097688523919$$



$$k_{14,0} = 91.632551202864647$$





- BEM widely used in homogeneous wave scattering applications (e.g. inverse obstacle problem) in the medium frequency range
- It is widely believed that there is no pollution in the boundary integral method but the proof is an open question
- This topic requires delicate analysis of oscillatory integrals (Lecture 2)
- If you want to treat non-smooth (Lipschitz) domains you need to use harmonic analysis techniques!