

The mass of super-Brownian motion upon exiting balls and Sheu's compact support condition

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Super-Brownian motion

Consider a finite-measure-valued strong Markov process $\{X_t : t \geq 0\}$ on \mathbb{R}^d whose evolution is characterised via its log-laplace semi-group: For all $f \in C_b^+(\mathbb{R})$, the space of positive, uniformly bounded, continuous functions on \mathbb{R}^d , and $\mu \in \mathcal{M}_F(\mathbb{R}^d)$ (the space of finite measure on \mathbb{R}^d),

$$-\log \mathbf{E}_\mu(e^{-\langle f, X_t \rangle}) = \int_{\mathbb{R}} v_f(x, t) \mu(dx), \quad t \geq 0,$$

where $v_f(x, t)$ is the unique positive solution to the evolution equation for $x \in \mathbb{R}$ and $t > 0$

$$\frac{\partial}{\partial t} v_f(x, t) = \frac{1}{2} \frac{\partial^2}{\partial x^2} v_f(x, t) - \psi(v_f(x, t)),$$

with initial condition $v_f(x, 0) = f(x)$. The branching mechanism ψ satisfies:

$$\psi(\lambda) = -\alpha\lambda + \beta\lambda^2 + \int_{(0, \infty)} (e^{-\lambda x} - 1 + \lambda x) \nu(dx), \quad (1)$$

for $\lambda \geq 0$ where $\alpha = -\psi'(0^+) \in (0, \infty)$, $\beta \geq 0$ and ν is a measure concentrated on $(0, \infty)$ which satisfies $\int_{(0, \infty)} (x \wedge x^2) \nu(dx) < \infty$.

Super-Brownian motion

- Another way of representing the log-Laplace semi-group evolution is via the integral equation:

$$v_f(x, t) = \mathbb{E}_x[f(\xi_t)] - \mathbb{E}_x \left[\int_0^t \psi(v_f(\xi_z, t - z)) dz \right],$$

and $((\xi_z, z \geq 0), \mathbb{P}_x)$ is an \mathbb{R}^d -Brownian motion with $\xi_0 = x$

- Choosing $f = 1$ produces the log-Laplace exponent a CSBP with branching mechanism ψ . That is to say the total mass process, $\|X_t\| := \langle 1, X_t \rangle$, $t \geq 0$, is a CSBP.
- This super-BM is the continuum analogue of Branching Brownian motion with a general off-spring distribution (including allowing for no offspring w.p.p.).
- The constant $-\psi'(0+) = \alpha$ gives us the growth rate and hence process is sub/super-critical with $\alpha < 0/\alpha > 0$. Largely indifferent to criticality in this talk.

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Mass decay and Grey's condition

- Following the same property of all (non-monotone) CSBPs:

$$\mathbf{P}_\mu(\lim_{t \rightarrow \infty} \|X_t\| = 0 \mid \|X_0\| = x) = e^{-\lambda^* \|\mu\|},$$

where $\psi(\lambda^*) = 0$ and $\mu \in \mathcal{M}_F(\mathbb{R}^d)$.

- On the event $\{\|X_t\| \rightarrow 0\}$, Grey (1974) gives us a nice dichotomy between the two ways in which this can happen: either

$\{\exists T(\omega) > 0$ s.t. $\|X_{T+t}\| = 0 \forall t \geq 0\}$ (extinction)

or $\{\|X_t\| \rightarrow 0$ and $\|X_t\| > 0 \forall t > 0\}$ (extinguishing)

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Sheu's condition

- For the case of super-Brownian motion, Sheu (1994) offers an additional unusual condition for the event of compact support: Let

$$\mathcal{S} = \overline{\bigcup_{t \geq 0} \text{supp} X_t}$$

Then for all compactly supported $\mu \in \mathcal{M}_F(\mathbb{R}^d)$,

$$\mathbf{P}_\mu(\mathcal{S} \text{ is compact}) = e^{-\lambda_* \|\mu\|}$$

if and only if

$$\int_0^\infty \frac{1}{\sqrt{\int_{\lambda_*}^\lambda \psi(\theta) \, d\theta}} \, d\lambda < \infty$$

and otherwise $\mathbb{P}_\mu(\mathcal{S} \text{ is compact}) = 0$.

- What is the relation between this condition and Grey's condition? Sheu's condition comes out of PDE analysis and it is unclear where the condition comes from.
- What is the relation between $\{\mathcal{S} \text{ is compact}\}$ and $\{\|X_t\| \rightarrow 0\}$?

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First passage branching process

- Fix an initial radius $r > 0$ and let $D_s := \{x \in \mathbb{R}^d : \|x\| < s\}$ be the open ball of radius $s \geq r$ around the origin.
- According to Dynkin's theory of exit measures we can describe the mass of X as it first exits the growing sequence of balls $(D_s, s \geq r)$ as a sequence of random measures on \mathbb{R}^d , known as branching Markov exit measures.
- We denote this sequence of branching Markov exit measures by $\{X_{D_s}, s \geq r\}$. Informally, the measure X_{D_s} is supported on the boundary ∂D_s and it is obtained by 'freezing' mass of the super-Brownian motion when it first hits ∂D_s . If X were a branching Brownian motion, then X_{D_s} would be a stopping line *à la* Chauvin-Neveu.
- For $s \geq r$, let $Z_s := \|X_{D_s}\|$ denote the mass that is 'frozen' when it first hits the boundary of the ball D_s . We can then define the mass process $(Z_s, s \geq r)$ which uses the radius s as its time-parameter.

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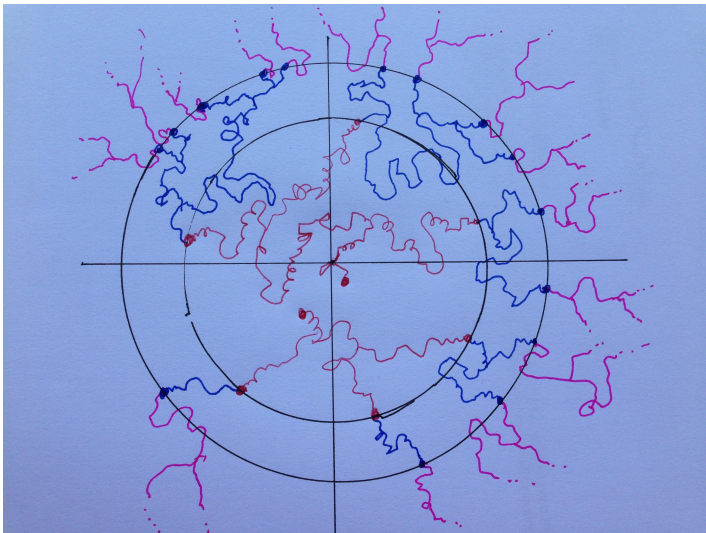
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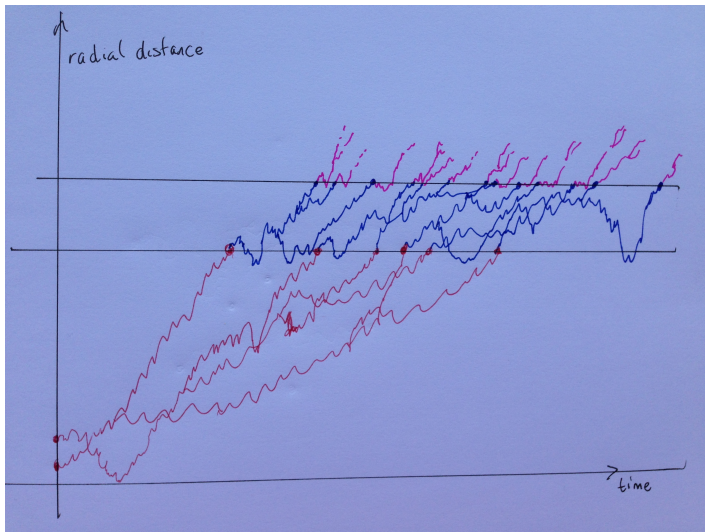
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Z is a time-inhomogenous CSBP



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Theorem

Let $r > 0$. The process $Z = (Z_s, s \geq r)$ is a time-inhomogeneous continuous-state branching process. Let $r > 0$ and $\mu \in \mathcal{M}_F(\partial D_r)$ with $\|\mu\| = a$. Then, for $s \geq r$, we have

$$E_{a,r}[e^{-\theta Z_s}] = e^{-u(r,s,\theta)a}, \quad \theta \geq 0,$$

where the Laplace functional $u(r, s, \theta)$ satisfies

$$u(r, s, \theta) = \theta - \int_r^s \Psi(z, u(z, s, \theta)) dz,$$

for a family of branching mechanisms $(\Psi(r, \cdot), r > 0)$ satisfying the PDE

$$\begin{aligned} \frac{\partial}{\partial r} \Psi(r, \theta) + \frac{1}{2} \frac{\partial}{\partial \theta} \Psi^2(r, \theta) + \frac{d-1}{r} \Psi(r, \theta) &= 2\psi(\theta) \\ \Psi(r, \lambda^*) &= 0, \end{aligned}$$

for $r > 0$, $\theta \in (0, \infty)$.

Asymptotic behaviour of Z

Proposition

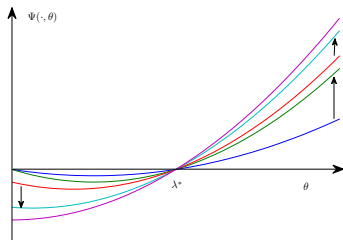
(i) For (sub)critical ψ , we have, for $0 < r \leq s$,

$$\Psi(r, \theta) \leq \Psi(s, \theta) \quad \text{for all } \theta \geq 0.$$

(ii) For supercritical ψ , we have, for $0 < r \leq s$,

$$\Psi(r, \theta) \geq \Psi(s, \theta) \quad \text{for all } \theta \leq \lambda^*$$

$$\Psi(r, \theta) \leq \Psi(s, \theta) \quad \text{for all } \theta \geq \lambda^*.$$



Asymptotic behaviour of Z

Lemma

For each $\theta \geq 0$, the limit $\lim_{r \uparrow \infty} \Psi(r, \theta) = \Psi_\infty(\theta)$ is finite and the convergence holds uniformly in θ on any bounded, closed subset of \mathbb{R}_+ . For any $\theta \geq 0$, we have

$$\Psi_\infty(\theta) = 2 \operatorname{sgn}(\psi(\theta)) \sqrt{\int_{\lambda^*}^{\theta} \psi(\lambda) \, d\lambda},$$

with $\lambda^* = 0$ in the (sub)critical case.

$$\begin{aligned} \frac{\partial}{\partial r} \Psi(r, \theta) + \frac{1}{2} \frac{\partial}{\partial \theta} \Psi^2(r, \theta) + \frac{d-1}{r} \Psi(r, \theta) &= 2\psi(\theta) \\ \Psi(r, \lambda^*) &= 0, \end{aligned}$$

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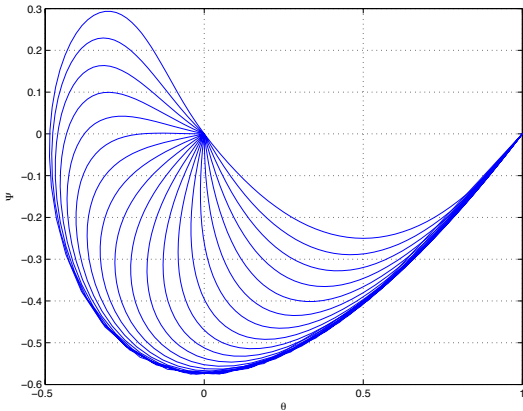
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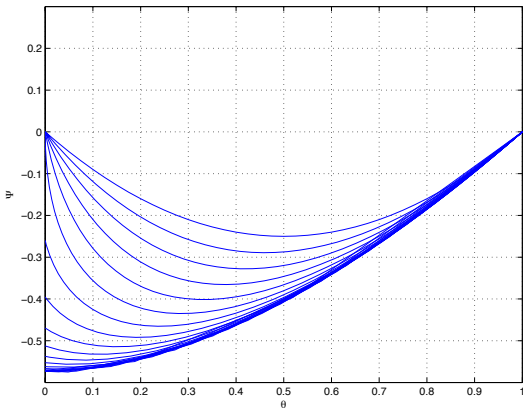
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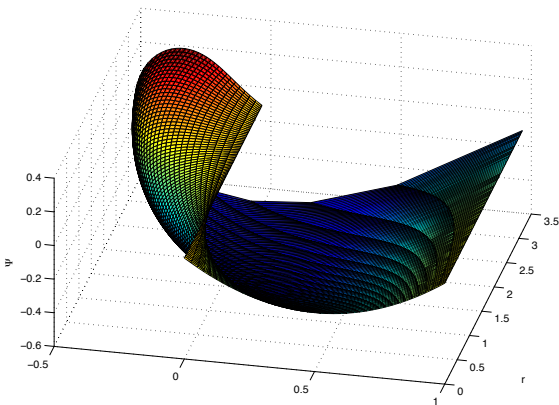
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Asymptotic behaviour of Z

Lemma

Denote by $((Z_s^\infty, s \geq 0), P^\infty)$ the standard CSBP associated with the limiting branching mechanism Ψ_∞ , with unit initial mass at time 0.

Then, for any $s > 0$, $\theta \geq 0$,

$$\lim_{r \rightarrow \infty} E_{r,1}[e^{-\theta Z_{r+s}}] = E^\infty[e^{-\theta Z_s^\infty}].$$

Sheu's condition is Grey's condition

- Sheu's condition is Grey's condition for Z^∞ .

$$\int^\infty \frac{1}{\sqrt{\int_{\lambda^*}^\lambda \psi(\theta) d\theta}} d\lambda = \int^\infty \frac{1}{\Psi^\infty(\lambda)} d\lambda.$$

- There is no hierarchy: $\{\|X_t\| \rightarrow 0\} \not\Leftarrow \{\mathcal{S} \text{ is compact}\}$.
- Take e.g. the supercritical branching mechanism $\psi(\lambda) = \lambda - (\lambda + 2)^\alpha + 2^\alpha$ for $\alpha \in (0, 1)$. This branching mechanism respects $\int^\infty 1/\psi(\lambda)d\lambda = \infty$ (extinguishing) but $\int^\infty 1/(\int_{\lambda^*}^\lambda \psi(\theta) d\theta)^{1/2}d\lambda = \infty$ (no compact support).
- In principle it could happen that ψ is such that we have a process that becomes extinct but which is not compactly supported.

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Martingales

- On the one hand, using the semi-group equations,

$$M_s^\lambda = e^{-\lambda^* Z_s} - \int_r^s \Psi(v, \lambda^*) Z_v e^{-\lambda^* Z_v} \mathbf{1}_{\{Z_v < \infty\}} dv, \quad s \geq r,$$

is a martingale.

- On the other hand $\exp\{-\lambda^* \|X_t\|\}$, $t \geq 0$ is a martingale since

$$\mathbf{E}_\mu \left[\mathbf{1}_{\{\|X_u\| \rightarrow 0\}} \mid \sigma(\|X_s\|, s \leq t) \right] = e^{-\lambda^* \|X_t\|}, \quad t \geq 0,$$

and hence so is

$$\mathbf{E}_\mu [\mathbf{1}_{\{\|X_v\| \rightarrow 0\}} \mid \sigma(\|X_{D_v}\|, r \leq v \leq s)] = e^{-\lambda^* \|X_{D_s}\|} = e^{-\lambda^* Z_s},$$

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Branching mechanism PDE

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- On the other hand: Recalling that the radial part of an \mathbb{R}^d -Brownian motion is a Bessel process, Dynkin's semigroup theory for branching Markov exit measures gives us

$$u(r, s, \theta) = \theta - \mathbb{E}_r^{\mathbb{R}} \int_0^{\tau_s} \psi(u(R_\ell, s, \theta)) d\ell, \quad 0 < r \leq s, \quad \theta \geq 0,$$

where $(R, \mathbb{P}^{\mathbb{R}})$ is a d -dimensional Bessel process and $\tau_s := \inf\{l > 0 : R_l > s\}$ its first passage time above level s

Branching mechanism PDE

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$$E_{a,r}[e^{-\theta Z_s}] = e^{-u(r,s,\theta)a}, \quad \theta \geq 0,$$

where the Laplace functional $u(r, s, \theta)$ satisfies

$$u(r, s, \theta) = \theta - \int_r^s \Psi(z, u(z, s, \theta)) dz,$$

- On the other hand: Recalling that the radial part of an \mathbb{R}^d -Brownian motion is a Bessel process, Dynkin's semigroup theory for branching Markov exit measures gives us

$$u(r, s, \theta) = \theta - \mathbb{E}_r^{\mathbb{R}} \int_0^{\tau_s} \psi(u(R_\ell, s, \theta)) d\ell, \quad 0 < r \leq s, \quad \theta \geq 0,$$

where $(R, \mathbb{P}^{\mathbb{R}})$ is a d -dimensional Bessel process and $\tau_s := \inf\{l > 0 : R_l > s\}$ its first passage time above level s

Branching mechanism PDE

- Define $\varphi(s) = \int_0^{r^2 s} R_\ell^{-2} d\ell$, $s \geq 0$, then

$$B_s = \log(r^{-1} R_{r^2 \varphi^{-1}(s)}), \quad s \geq 0,$$

is a one-dimensional Brownian motion with drift $\frac{d}{2} - 1$.

- the Last semi-group equation can be developed into

$$\begin{aligned} u(r, s, \theta) &= \theta - \mathbb{E}_{\log r} \int_0^{T_{\log s}} \psi(u(e^{B_\ell}, s, \theta)) e^{2B_\ell} d\ell \\ &= \mathbb{E}_{\log r} \sum_{\log r \leq u \leq \log s} \int_0^{\zeta(u)} \psi(u(e^{u-e_u(\ell)}, s, \theta)) e^{2(u-e_u(\ell))} d\ell \\ &= \theta - 2 \int_r^s v^{1-d} \int_0^v \psi(u(z, s, \theta)) z^{d-1} dz dv. \end{aligned}$$

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- Line up the two representations of the semi-group equation:

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and

$$u(r, s, \theta) = \theta - \int_r^s \Psi(z, u(z, s, \theta)) dz,$$

- Fiddling with derivatives in s, r and θ , gives the desired PDE

$$\begin{aligned} \frac{\partial}{\partial r} \Psi(r, \theta) + \frac{1}{2} \frac{\partial}{\partial \theta} \Psi^2(r, \theta) + \frac{d-1}{r} \Psi(r, \theta) &= 2\psi(\theta) \\ \Psi(r, \lambda^*) &= 0, \end{aligned}$$

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