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<sup>&</sup>lt;sup>1</sup>Joint work with Terence Chan and Mladen Savov

Analytical properties of scale functions for spectrally negative Lévy processes

## Spectrally negative Lévy processes and scale functions

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- $\blacksquare$  Let us define the Laplace exponent  $\psi$  on  $[0,\infty)$  by

$$\mathbb{E}_0(e^{\beta X_t}) = e^{\psi(\beta)t}.$$

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Then for  $q \ge 0$  we may define the *q*-scale function  $W^{(q)} : \mathbb{R} \to [0,\infty)$  by  $W^{(q)}(x) = 0$  for x < 0 and on  $(0,\infty)$  it is the unique right continuous function such that for  $\beta > \Phi(q)$ 

$$\int_0^\infty e^{-\beta x} W^{(q)}(x) dx = \frac{1}{\psi(\beta) - q}$$

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- Note, formally we need to prove that scale functions exist they do!

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Scale functions

## Why are scale functions important?

 They appear in virtually all fluctuation identities for spectrally negative Lévy processes.

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- For example, resolvent in a strip: for any a > 0,  $x, y \in [0, a]$ ,  $q \ge 0$

$$\int_0^\infty e^{-qt} \mathbb{P}_x(X_t \in dy, \ t < \tau_a^+ \land \tau_0^-) dt$$
$$= \left\{ \frac{W^{(q)}(x) W^{(q)}(a-y)}{W^{(q)}(a)} - W^{(q)}(x-y) \right\} dy.$$

where

$$\tau_a^+ = \inf\{t > 0 : X_t > a\}$$
 and  $\tau_0^- = \inf\{t > 0 : X_t < 0\}.$ 

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where

$$\tau_a^+ = \inf\{t > 0 : X_t > a\} \text{ and } \tau_0^- = \inf\{t > 0 : X_t < 0\}.$$

• Or another example is: for a > 0,  $x \in [0, a]$ ,

$$\mathbb{E}_{x}(e^{-q\tau_{0}^{-}}\mathbf{1}_{\{\tau_{0}^{-}<\tau_{a}^{+}\}}) = Z^{(q)}(x) - W^{(q)}(x)\frac{Z^{(q)}(a)}{W^{(q)}(a)}$$

where

$$Z^{(q)}(x) = 1 + q \int_0^x W^{(q)}(y) dy.$$

## A 'basis' of fundamental solutions to the 'generator equation'

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It is not difficult to check that for  $t \ge 0$  and  $a \in (0,\infty]$ 

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are martingales.

Suppose that Γ is the infinitessimal generator of X. Then another way of expressing these martingale properties is by writing (in a loose sense)

$$(\Gamma - q) W^{(q)}(x) = 0$$
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• As many known fluctuation identites turn out to be linear combinations of  $W^{(q)}$  and  $Z^{(q)}$ , this suggest that potentially a small 'theory' would be possible in which these functions play the role of a fundamental basis of solutions to the equation

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Before even pursuing that objective, one needs to ask whether  $(\Gamma - q) W^{(q)}(x)$  makes sense mathematically.

# How smooth is $W^{(q)}$ ? Where to start looking?

-Scale functions

## How smooth is $W^{(q)}$ ? Where to start looking?

• Enough to answer the question for q = 0 and  $\mathbb{E}(X_1) = \psi'(0+) \ge 0$  as all other cases can be reduced to this case via the relation

$$W^{(q)}(x) = e^{\Phi(q)x} W_{\Phi(q)}(x)$$

where  $W_{\Phi(q)}$  plays the role of W after an exponential change of measure under which X drifts to  $+\infty$ .

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  - Excursion theory

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- Two key theories that will help analyse the issue of smoothness:
  - Excursion theory
  - Renewal theory

## The connection with excursion theory

#### The connection with excursion theory

It is known that

$$W(x) = W(a) \exp\left\{-\int_{x}^{a} \nu(\overline{\epsilon} > t) dt\right\}$$

where  $\nu(\cdot)$  is the intensity measure associated with the Poisson point process of excursions and  $\epsilon$  is the canonical excursion and  $\overline{\epsilon}$  is its maximum value.

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Hence

$$W'(x+) - W'(x-) = \nu(\overline{\epsilon} = x)$$

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(which immediately implies that for unbounded variation processes  $W \in C^1(0,\infty)$ ).

## The connection with excursion theory

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#### The connection with excursion theory

For processes with bounded variation paths (then necessarily  $X_t = \delta t - S_t$  where  $\delta > 0$  and S is a driftless subordinator):

$$\nu(\overline{\epsilon} > t) = \frac{1}{\delta} \Pi(-\infty, -t) + \frac{1}{\delta} \int_{[-t,0]} \Pi(dz) \left(1 - \frac{W(t+z)}{W(t)}\right)$$

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showing  $W \in C^1(0,\infty) \Leftrightarrow \Pi$  has no atoms.

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showing  $W \in C^1(0,\infty) \Leftrightarrow \Pi$  has no atoms.

In the presence of a Gaussian component ( $\sigma \neq 0$ ):

$$\nu(\epsilon \text{ passes } x \text{ continuously}) = \frac{\sigma^2}{2} \left( \frac{W'(x)^2}{W(x)} - W''(x) \right)$$

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showing that  $W \in C^2(0,\infty)$ .

## The connection with renewal theory

Scale functions

#### The connection with renewal theory

An important observation which helps us understand what kind of answer we might expect comes from the Wiener-Hopf factorization:

$$\psi(\beta) = \beta \phi(\beta)$$

where  $\phi$  is the Laplace exponent of the descending ladder height process and hence an integration by parts gives

$$\int_{[0,\infty)} e^{-\beta x} W(dx) = \frac{1}{\phi(\beta)}.$$

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This means that W can be seen as the renewal function of a (killed) subordinator, say H, which has Laplace exponent  $\phi$ , from which it can be calculated that its jump measure is given by  $\Pi(-\infty, -x)dx$ , its drift is given  $\sigma^2/2$  and its killing rate  $\mathbb{E}(X_1) > 0$ . Here, renewal function means

$$W(dx) = \int_0^\infty \mathbb{P}(H_t \in dx) dt$$

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## The connection with renewal theory

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Then appealing to the formula

$$1 = \mathbb{P}(H_{\tau^+_x} = x) + \mathbb{P}(H_{\tau^+_x} > x)$$

using Kesten's classical result for the probability of continuous crossing for a subordinator and Kesten-Horowitz-Bertoin formula for overshoots of subordinators one derrives that

$$1 = \frac{\sigma^2}{2} W'(x) + \int_0^x W'(x-y) \{\overline{\overline{\Pi}}(y) + \mathbb{E}(X_1)\} dy$$

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This can be seen as a renewal equation of the form f = 1 + f \* g for appropriate f and g. But only when  $\sigma \neq 0$ , otherwise it takes the form f = f \* g.

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- This can be seen as a renewal equation of the form f = 1 + f \* g for appropriate f and g. But only when  $\sigma \neq 0$ , otherwise it takes the form f = f \* g.
- When X has bounded variation paths  $(X_t = \delta t S_t)$  it a direct inverse of the Laplace transform for W also gives us

$$\delta W(x) = 1 + \int_0^x W(x - y) \Pi(-\infty, -y) dy$$

which is again a renewal function of the form f = 1 + f \* g.

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## Some more results

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 $\blacksquare$  Suppose that  $\sigma \neq 0$  and

$$\inf\{\beta \ge 0: \int_{|x|<1} |x|^{\beta} \Pi(dx) < \infty\} \in [0,2).$$

Then for  $k = 0, 1, 2, \cdots$   $W \in C^{k+3}(0, \infty)$  if and only if  $\overline{\Pi} \in C^k(0, \infty)$ .

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Suppose that X has paths of bounded variation and  $-\overline{\Pi}$  has a derivative  $\pi$  such that  $\pi(x) \leq Cx^{-1-\alpha}$  in the neighbourhood of the origin for some  $\alpha < 1$  and C > 0. Then for  $k = 0, 1, 2, \cdots$   $W \in C^{k+3}(0, \infty)$  if and only if  $\overline{\Pi} \in C^k(0, \infty)$ .

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## Doney's conjecture for scale functions

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• If  $\sigma \neq 0$  (X has a Gaussian component) then

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