Skeletal stochastic differential equations for continuous-state branching process

A. Kyprianou, D. Fekete, J. Fontbona



DEFINITION OF ψ -CSBP.

A CSBP (X, \mathbb{P}_x) is a non-negative valued strong Markov process with probabilities $(\mathbb{P}_x, x \ge 0)$ such that for any $x, y \ge 0$, $\mathbb{P}_{x+y} = \mathbb{P}_x * \mathbb{P}_y$.

In particular

$$\mathbb{E}_{x}(\mathrm{e}^{-\theta X_{t}})=\mathrm{e}^{-xu_{t}(\theta)}, \qquad x,\theta,t\geq 0,$$

where $u_t(\theta)$ uniquely solves the evolution equation

$$u_t(\theta) + \int_0^t \psi(u_s(\theta)) \mathrm{d}s = \theta, \qquad t \ge 0.$$

Here, we assume that the so-called branching mechanism ψ takes the form

$$\psi(\theta) = -\alpha\theta + \beta\theta^2 + \int_{(0,\infty)} (e^{-\theta x} - 1 + \theta x) \Pi(dx), \ \theta \ge 0,$$

where $\alpha \in \mathbb{R}$, $\beta \geq 0$ and Π is a measure concentrated on $(0, \infty)$ which satisfies $\int_{(0,\infty)} (x \wedge x^2) \Pi(dx) < \infty$.

PROPERTIES.

We assume that the process is conservative, i.e.

$$\int_{0+} \frac{1}{|\psi(\xi)|} \mathrm{d}\xi = \infty.$$

It is easily verified that

$$\mathbb{E}_{x}[X_{t}] = x \mathrm{e}^{-\psi'(0+)t}, \qquad t, x \ge 0.$$

We say that the CSBP is supercritical, critical or subcritical accordingly as $-\psi'(0+) = \alpha$ is strictly positive, equal to zero or strictly negative.

For a **supercritical** ψ -CSBP the probability of extinction is

$$\mathbb{P}_x(\lim_{t\uparrow\infty}X_t=0)=\mathrm{e}^{-\lambda^*x},$$

where λ^* is the unique root on $(0, \infty)$ of the equation $\psi(\theta) = 0$.



PROLIFIC SKELETON I.

The supercritical ψ -CSBP is equal in law to the total mass process obtained by the following construction.

▶ Initiate $Po(\lambda^* x)$ independent Galton-Watson processes with branching generator

$$q\left(\sum_{k\geq 0} p_k r^k - r\right) = \frac{1}{\lambda^*} \psi(\lambda^*(1-r)), \qquad r \in [0,1],$$

where $q = \psi'(\lambda^*)$, $p_0 = p_1 = 0$ and for $k \ge 2$

$$p_k = \frac{1}{\lambda^* \psi'(\lambda^*)} \left\{ \beta(\lambda^*)^2 \mathbf{1}_{\{k=2\}} + (\lambda^*)^k \int_{(0,\infty)} \frac{r^k}{k!} e^{-\lambda^* r} \Pi(dr) \right\}.$$

Along the edges immigrate CSBPs at rate

$$2\beta d\mathbb{Q}^* + \int_0^\infty y \mathrm{e}^{-\lambda^* y} \Pi(\mathrm{d} y) \mathrm{d} \mathbb{P}_y^*,$$

where \mathbb{P}_{x}^{*} , $x \ge 0$ is the law of the CSBP with branching mechanism $\psi^{*}(\lambda) = \psi(\lambda + \lambda^{*})$ and \mathbb{Q}^{*} is the associated excursion measure.

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PROLIFIC SKELETON II.

Given that an individual dies and branches into $k \ge 2$ offspring, an independent ψ^* -CSBP is immigrated with initial mass r with probability

$$\eta_k(\mathrm{d}r) = \frac{1}{p_k \lambda^* \psi'(\lambda^*)} \left\{ \beta(\lambda^*)^2 \delta_0(\mathrm{d}r) \mathbf{1}_{\{k=2\}} + (\lambda^*)^k \frac{r^k}{k!} \mathrm{e}^{-\lambda^* r} \Pi(\mathrm{d}r) \right\}.$$

Finally an independent ψ^* -CSBP is issued at time zero with initial mass x.



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λ -skeleton I.

Let $\lambda \geq \lambda^*$. Define the Esscher transformed branching mechanism $\psi_{\lambda} : \mathbb{R}_+ \to \mathbb{R}_+$ for $\theta \geq -\lambda$ and $\lambda \geq \lambda^*$ by $\psi_{\lambda}(\theta) = \psi(\theta + \lambda) - \psi(\lambda)$.

The supercritical ψ -CSBP is equal in law to the total mass process obtained by the following construction.

• Initiate $Po(\lambda x)$ independent Galton-Watson processes with branching generator

$$q\left(\sum_{k\geq 0}p_kr^k-r\right)=\frac{1}{\lambda}\psi(\lambda(1-r)),\qquad r\in[0,1],$$

where $q = \psi'(\lambda)$, $p_0 = \psi(\lambda) / \lambda \psi'(\lambda)$, $p_1 = 0$ and for $k \ge 2$

$$p_k = \frac{1}{\lambda \psi'(\lambda)} \left\{ \beta \lambda^2 \mathbf{1}_{\{k=2\}} + \int_{(0,\infty)} \frac{(\lambda r)^k}{k!} \mathrm{e}^{-\lambda r} \Pi(\mathrm{d}r) \right\}.$$

λ -skeleton II.

Along the edges immigrate CSBPs at rate

$$2\beta d\mathbb{Q}^{(\lambda)} + \int_0^\infty y \mathrm{e}^{-\lambda y} \Pi(\mathrm{d} y) \mathrm{d} \mathbb{P}_y^{(\lambda)},$$

where $\mathbb{P}_x^{(\lambda)}$, $x \ge 0$ is the law of the CSBP with branching mechanism ψ_{λ} and $\mathbb{Q}^{(\lambda)}$ is the associated excursion measure.

Given that an individual dies and branches into $k \in \mathbb{N}_0 \setminus \{1\}$ offspring, an independent ψ_{λ} -CSBP is immigrated with initial mass r with probability

$$\begin{split} \eta_k(\mathrm{d}r) &= \frac{1}{p_k \lambda \psi'(\lambda)} \left\{ \psi(\lambda) \mathbf{1}_{\{k=0\}} \delta_0(\mathrm{d}r) + \beta \lambda^2 \mathbf{1}_{\{k=2\}} \delta_0(\mathrm{d}r) \right. \\ &\left. + \mathbf{1}_{\{k\geq 2\}} \frac{(\lambda r)^k}{k!} \mathrm{e}^{-\lambda r} \Pi(\mathrm{d}r) \right\}, \end{split}$$

Finally an independent ψ_{λ} -CSBP is issued at time zero with initial mass *x*.

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Introduction	Supercritical CSBP	Subcritical CSBP
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SDE.

The process (X, \mathbb{P}_x) , x > 0, can be represented as the unique strong solution to the stochastic differential equation (SDE)

$$X_{t} = x + \alpha \int_{0}^{t} X_{s-} ds + \sqrt{2\beta} \int_{0}^{t} \int_{0}^{X_{s-}} W(ds, du) + \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{X_{s-}} r \tilde{N}(ds, dr, d\nu),$$
(1)

for $x > 0, t \ge 0$, where

- ▶ W(ds, du) is a white noise process on $(0, \infty)^2$ based on the Lebesgue measure $ds \otimes du$,
- ▶ $N(ds, dr, d\nu)$ is a Poisson point process on $(0, \infty)^3$ with intensity $ds \otimes \Pi(dr) \otimes d\nu$, and $\tilde{N}(ds, dr, d\nu)$ the compensated measure of $N(ds, dr, d\nu)$.

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THINNING OF THE SDE I.

We can introduce an additional mark to atoms of *N*, resulting in an 'extended' Poisson random measure, $\mathcal{N}(ds, dr, d\nu, dk)$ on $(0, \infty)^3 \times \mathbb{N}_0$ with intensity

$$\mathrm{d} s\otimes \Pi(\mathrm{d} r)\otimes \mathrm{d} \nu\otimes rac{(\lambda r)^k}{k!}\mathrm{e}^{-\lambda r}\sharp(\mathrm{d} k).$$

Define three random measures by

$$\begin{split} N^0(\mathrm{d} s,\mathrm{d} r,\mathrm{d} \nu) &= \mathcal{N}(\mathrm{d} s,\mathrm{d} r,\mathrm{d} \nu,\{k=0\}),\\ N^1(\mathrm{d} s,\mathrm{d} r,\mathrm{d} \nu) &= \mathcal{N}(\mathrm{d} s,\mathrm{d} r,\mathrm{d} \nu,\{k=1\}) \end{split}$$

and

$$N^{2}(\mathrm{d} s, \mathrm{d} r, \mathrm{d} \nu) = \mathcal{N}(\mathrm{d} s, \mathrm{d} r, \mathrm{d} \nu, \{k \geq 2\}).$$

We have that N^0 , N^1 and N^2 are independent Poisson point processes on $(0, \infty)^3$ with respective intensities $ds \otimes e^{-\lambda r} \Pi(dr) \otimes d\nu$, $ds \otimes (\lambda r) e^{-\lambda r} \Pi(dr) \otimes d\nu$ and $ds \otimes \sum_{k=2}^{\infty} (\lambda r)^k e^{-\lambda r} \Pi(dr) / k! \otimes d\nu$.

THINNING OF THE SDE II.

$$\begin{split} X_t &= x + \alpha \int_0^t X_{s-} ds + \sqrt{2\beta} \int_0^t \int_0^{X_{s-}} W(ds, du) + \int_0^t \int_0^\infty \int_0^{X_{s-}} r \tilde{N}^0(ds, dr, d\nu) \\ &+ \int_0^t \int_0^\infty \int_0^{X_{s-}} r N^1(ds, dr, d\nu) + \int_0^t \int_0^\infty \int_0^{X_{s-}} r N^2(ds, dr, d\nu) \\ &- \int_0^t \int_0^\infty X_{s-} \sum_{n=1}^\infty \frac{(\lambda r)^n}{n!} e^{-\lambda r} r \Pi(dr) ds \\ &= x - \psi'(\lambda) \int_0^t X_s ds + \sqrt{2\beta} \int_0^t \int_0^{X_{s-}} W(ds, du) + \int_0^t \int_0^\infty \int_0^{X_{s-}} r \tilde{N}^0(ds, dr, d\nu) \\ &+ \int_0^t \int_0^\infty \int_0^{X_{s-}} r N^1(ds, dr, d\nu) + 2\beta\lambda \int_0^t X_{s-} ds \\ &+ \int_0^t \int_0^\infty \int_0^{X_{s-}} r N^2(ds, dr, d\nu), \end{split}$$

(In the last equality we have used that $-\int_{(0,\infty)} (1-e^{-\lambda r})r\Pi(dr) = -\alpha + 2\beta\lambda - \psi'(\lambda)$).

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THINNING OF THE SDE II.

$$\begin{split} \mathbf{X}_{t} &= x + \alpha \int_{0}^{t} X_{s-} \mathrm{ds} + \sqrt{2\beta} \int_{0}^{t} \int_{0}^{X_{s-}} W(\mathrm{ds}, \mathrm{d}u) + \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{X_{s-}} r \tilde{N}^{0}(\mathrm{ds}, \mathrm{d}r, \mathrm{d}\nu) \\ &+ \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{X_{s-}} r N^{1}(\mathrm{ds}, \mathrm{d}r, \mathrm{d}\nu) + \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{X_{s-}} r N^{2}(\mathrm{ds}, \mathrm{d}r, \mathrm{d}\nu) \\ &- \int_{0}^{t} \int_{0}^{\infty} X_{s-} \sum_{n=1}^{\infty} \frac{(\lambda r)^{n}}{n!} \mathrm{e}^{-\lambda r} r \Pi(\mathrm{d}r) \mathrm{ds} \\ &= x - \psi'(\lambda) \int_{0}^{t} X_{s} \mathrm{ds} + \sqrt{2\beta} \int_{0}^{t} \int_{0}^{X_{s-}} W(\mathrm{ds}, \mathrm{d}u) + \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{X_{s-}} r \tilde{N}^{0}(\mathrm{ds}, \mathrm{d}r, \mathrm{d}\nu) \\ &+ \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{X_{s-}} r N^{1}(\mathrm{ds}, \mathrm{d}r, \mathrm{d}\nu) + 2\beta\lambda \int_{0}^{t} X_{s-} \mathrm{ds} \\ &+ \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{X_{s-}} r N^{2}(\mathrm{ds}, \mathrm{d}r, \mathrm{d}\nu), \end{split}$$

(In the last equality we have used that $-\int_{(0,\infty)} (1-e^{-\lambda r})r\Pi(dr) = -\alpha + 2\beta\lambda - \psi'(\lambda)$).

Theorem

Suppose that ψ corresponds to a supercritical branching mechanism (i.e. $\alpha > 0$) and $\lambda \ge \lambda^*$. Consider the coupled system of SDEs

Supercritical CSBP

$$\begin{pmatrix} \Lambda_t \\ Z_t \end{pmatrix} = \begin{pmatrix} \Lambda_0 \\ Z_0 \end{pmatrix} - \psi'(\lambda) \int_0^t \begin{pmatrix} \Lambda_{s-} \\ 0 \end{pmatrix} ds + \sqrt{2\beta} \int_0^t \int_0^{\Lambda_{s-}} \begin{pmatrix} 1 \\ 0 \end{pmatrix} W(ds, du) + \int_0^t \int_0^\infty \int_0^{\Lambda_{s-}} \begin{pmatrix} r \\ 0 \end{pmatrix} \tilde{N}^0(ds, dr, d\nu) + \int_0^t \int_0^\infty \int_1^{Z_{s-}} \begin{pmatrix} r \\ 0 \end{pmatrix} N^1(ds, dr, dj) + \int_0^t \int_0^\infty \int_0^\infty \int_1^{Z_{s-}} \begin{pmatrix} r \\ k-1 \end{pmatrix} N^2(ds, dr, dk, dj) + 2\beta \int_0^t \begin{pmatrix} Z_{s-} \\ 0 \end{pmatrix} ds, \quad t \ge 0,$$
(2)

with $\Lambda_0 \ge 0$ given and fixed. Under the assumption that Z_0 is an independent random variable which is Poisson distributed with intensity $\lambda \Lambda_0$ the system (2) has a unique strong solution such that:

- (i) For t ≥ 0, Z_t | F_t^Λ is Poisson distributed with intensity λΛ_t, where F_t^Λ := σ(Λ_s : s ≤ t);
- (ii) The process $(\Lambda_t, t \ge 0)$ is a weak solution to (1).

(DRIVING SOURCES OF RANDOMNESS I.)

Let
$$\mathbb{N}_0 = \{0\} \cup \mathbb{N}$$
 and $\sharp(d\ell) = \sum_{i \in \mathbb{N}_0} \delta_i(d\ell), \ell \ge 0$.

Then in the previous theorem

- ▶ \mathbb{N}^0 is a Poisson random measure on $(0, \infty)^3$ with intensity measure $ds \otimes e^{-\lambda r} \Pi(dr) \otimes d\nu$, $\tilde{\mathbb{N}}^0$ is the associated compensated version of \mathbb{N}^0 ,
- ▶ $\mathbb{N}^1(ds, dr, dj)$ is a Poisson point process on $(0, \infty)^2 \times \mathbb{N}$ with intensity $ds \otimes re^{-\lambda r} \Pi(dr) \otimes \sharp(dj)$,
- ▶ $\mathbb{N}^2(ds, dr, dk, dj)$ is a Poisson point process on $(0, \infty)^2 \times \mathbb{N}_0 \times \mathbb{N}$ with intensity $\psi'(\lambda) ds \otimes \eta_k(dr) \otimes p_k \sharp(dk) \otimes \sharp(dj)$, and
- ▶ W(ds, du) is the white noise process on $(0, \infty)^2$ based on the Lebesgue measure $ds \otimes du$.

SUBCRITICAL CSBP.



RAY-KNIGHT REPRESENTATION.

Assume that Grey's condition is satisfied, ie.

$$\int^{\infty} \frac{1}{\psi(\theta)} \mathrm{d}u < \infty.$$

Let

► $(\xi_t, t \ge 0)$ be a spectrally positive Lévy process with Laplace exponent ψ ,

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► $(\hat{\xi}_r^{(t)}, 0 \le r \le t)$, where $\hat{\xi}_r^{(t)} := \xi_t - \xi_{(t-r)-}$, the time reversed process at time *t*,

$$\blacktriangleright \hat{S}_r^{(t)} := \sup_{s \le r} \hat{\xi}_s^{(t)}.$$

The process $(H_t, t \ge 0)$ is called the height process if H_t is the local time at level 0, at time *t* of $\hat{S}^{(t)} - \hat{\xi}^{(t)}$. Denote by L_t^a the local time up to time *t* of *H* at level $a \ge 0$, and let $T_x := \inf\{t \ge 0 : \xi_t = -x\}$.

Then the generalised Ray-Knight theorem for the ψ -CSBP process states that $(L^a_{T_v}, a \ge 0)$ has a càdàg modification for which

$$(L_{T_x}^t, t \ge 0) \stackrel{d}{=} (X, \mathbb{P}_x),$$

that is, the two processes are equal in law.

GENEALOGY OF SUBCRITICAL CSBP.

Excursions of *H* away from 0 form a PPP, denote by n its intensity, and let ϵ be a canonical excursion under n.

Let $\zeta = \inf\{s > 0, \epsilon_s = 0\}$, and define

$$d_{\epsilon}(s,t) = \epsilon_s + \epsilon_t - \inf_{s \wedge t \le r \le s \lor t} \epsilon_r, \quad (s,t) \in [0,\zeta]^2.$$

Then we can define the equivalence relation \sim_{ϵ} , such that $(s \sim_{\epsilon} t)$ is and only if $d_{\epsilon}(s,t) = 0$, and $\mathcal{T}_{\epsilon} = [0,\zeta] \setminus \sim_{\epsilon}$.

The compact metric space $(\mathcal{T}_{\epsilon}, d_{\epsilon})$ is called a Lévy random tree.

height



Fix T > 0. Define $(Z_t^T, 0 \le t < T)$ as the process that counts the number of excursions above level *t* that hit level *T*.

Then Z^T is a time-dependent continuous-time Galton-Watson process which at time t branching at rate

$$q^{T-t} = \frac{u_{T-t}(\infty)\psi'(u_{T-t}(\infty)) - \psi(u_{T-t}(\infty))}{u_{T-t}(\infty)}, \qquad t \in [0,T),$$

and its offspring distribution $(p_k^{T-t}, k \ge 0)$ is given by $p_0^{T-t} = p_1^{T-t} = 0$,

$$p_k^{T-t} = \frac{1}{u_{T-t}(\infty)q^{T-t}} \times \left\{ \beta u_{T-t}^2(\infty) \mathbf{1}_{\{k=2\}} + \int_0^\infty \frac{(u_{T-t}(\infty)x)^k}{k!} \mathrm{e}^{-u_{T-t}(\infty)x} \Pi(dx) \right\}.$$



Fix T > 0. Define $(Z_t^T, 0 \le t < T)$ as the process that counts the number of excursions above level *t* that hit level *T*.

Then Z^T is a time-dependent continuous-time Galton-Watson process which at time t branching at rate

$$q^{T-t} = \frac{u_{T-t}(\infty)\psi'(u_{T-t}(\infty)) - \psi(u_{T-t}(\infty))}{u_{T-t}(\infty)}, \qquad t \in [0,T),$$

and its offspring distribution $(p_k^{T-t}, k \ge 0)$ is given by $p_0^{T-t} = p_1^{T-t} = 0$,

$$p_k^{T-t} = \frac{1}{u_{T-t}(\infty)q^{T-t}} \times \left\{ \beta u_{T-t}^2(\infty) \mathbf{1}_{\{k=2\}} + \int_0^\infty \frac{(u_{T-t}(\infty)x)^k}{k!} \mathrm{e}^{-u_{T-t}(\infty)x} \Pi(dx) \right\}.$$



Fix T > 0. Define $(Z_t^T, 0 \le t < T)$ as the process that counts the number of excursions above level t that hit level T.

Then Z^T is a time-dependent continuous-time Galton-Watson process which at time t branching at rate

$$q^{T-t} = \frac{u_{T-t}(\infty)\psi'(u_{T-t}(\infty)) - \psi(u_{T-t}(\infty))}{u_{T-t}(\infty)}, \qquad t \in [0,T),$$

and its offspring distribution $(p_k^{T-t}, k \ge 0)$ is given by $p_0^{T-t} = p_1^{T-t} = 0$,

$$p_k^{T-t} = \frac{1}{u_{T-t}(\infty)q^{T-t}} \times \left\{ \beta u_{T-t}^2(\infty) \mathbf{1}_{\{k=2\}} + \int_0^\infty \frac{(u_{T-t}(\infty)x)^k}{k!} \mathrm{e}^{-u_{T-t}(\infty)x} \Pi(dx) \right\}.$$



Fix T > 0. Define $(Z_t^T, 0 \le t < T)$ as the process that counts the number of excursions above level *t* that hit level *T*.

Then Z^T is a time-dependent continuous-time Galton-Watson process which at time t branching at rate

$$q^{T-t} = \frac{u_{T-t}(\infty)\psi'(u_{T-t}(\infty)) - \psi(u_{T-t}(\infty))}{u_{T-t}(\infty)}, \qquad t \in [0,T),$$

and its offspring distribution $(p_k^{T-t}, k \ge 0)$ is given by $p_0^{T-t} = p_1^{T-t} = 0$,

$$p_k^{T-t} = \frac{1}{u_{T-t}(\infty)q^{T-t}} \times \left\{ \beta u_{T-t}^2(\infty) \mathbf{1}_{\{k=2\}} + \int_0^\infty \frac{(u_{T-t}(\infty)x)^k}{k!} \mathrm{e}^{-u_{T-t}(\infty)x} \Pi(dx) \right\}.$$



IMMIGRATION.

As

$$\mathbb{P}\left[X_T=0|\mathcal{F}_t\right]=\mathrm{e}^{-X_tu_{T-t}(\infty)},$$

the law of X conditioned to die out by time T can be obtained by the following change of measure

$$\frac{\mathrm{d}\mathbb{P}_x^T}{\mathrm{d}\mathbb{P}_x}\Big|_{\mathcal{F}_t} = \frac{\mathrm{e}^{-X_t u_{T-t}(\infty)}}{\mathrm{e}^{-x u_T(\infty)}}, \qquad t \ge 0, x > 0.$$

We get that (X, \mathbb{P}_x^T) is a time-dependent CSBP with Laplace transform

$$\mathbb{E}_x^T[\mathrm{e}^{-\theta X_t}] = \mathrm{e}^{-xV_t^T(\theta)}, \quad 0 \le t < T, \ x, \theta \ge 0,$$

where

$$V_t^T(\theta) = u_t(\theta + u_{T-t}(\infty)) - u_T(\infty).$$

Note that

$$\lim_{t\to T} u_{T-t}(\infty) = \infty, \text{ and } \lim_{T\to\infty} u_{T-t}(\infty) = 0.$$

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Theorem

Suppose that ψ corresponds to a (sub)critical branching mechanism (i.e. $\alpha \leq 0$) which satisfies Grey's condition. Fix a time horizon T > 0 and consider the coupled system of SDEs

$$\begin{pmatrix} \Lambda_t^T \\ Z_t^T \end{pmatrix} = \begin{pmatrix} \Lambda_0^T \\ Z_0^T \end{pmatrix} - \int_0^t \psi'(u_{T-s}(\infty)) \begin{pmatrix} \Lambda_{s-}^T \\ 0 \end{pmatrix} ds + \sqrt{2\beta} \int_0^t \int_0^{\Lambda_{s-}^I} \begin{pmatrix} 1 \\ 0 \end{pmatrix} W(ds, du)$$

$$+ \int_0^t \int_0^\infty \int_0^{\Lambda_{s-}^T} \begin{pmatrix} r \\ 0 \end{pmatrix} \tilde{N}_T^0(ds, dr, d\nu)$$

$$+ \int_0^t \int_0^\infty \int_1^{Z_{s-}^T} \begin{pmatrix} r \\ 0 \end{pmatrix} N_T^1(ds, dr, dj)$$

$$+ \int_0^t \int_0^\infty \int_0^\infty \int_1^{Z_{s-}^T} \begin{pmatrix} r \\ k-1 \end{pmatrix} N_T^2(ds, dr, dk, dj)$$

$$+ 2\beta \int_0^t \begin{pmatrix} Z_{s-}^T \\ 0 \end{pmatrix} ds, \quad 0 \le t < T.$$

$$(3)$$

with $\Lambda_0^T \ge 0$ given and fixed. Under the assumption that Z_0^T is an independent random variable which is Poisson distributed with intensity $u_T(\infty)\Lambda_0^T$ the system (3) has a unique strong solution such that:

- (i) For $T > t \ge 0$, $Z_t^T | \mathcal{F}_t^{\Lambda^T}$ is Poisson distributed with intensity $u_{T-t}(\infty) \Lambda_t^T$, where $\mathcal{F}_t^{\Lambda^T} := \sigma(\Lambda_s^T : s \le t)$;
- (ii) Conditional on $(\mathcal{F}^{\Lambda_t^T}, 0 \le t < T)$, the process $(\Lambda_{t=0}^T, 0 \le t < T)$ is a weak

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(DRIVING FORCES OF RANDOMNESS II.)

In the previous theorem

- ▶ \mathbb{N}_T^0 is a Poisson random measure on $[0, \infty)^3$ with intensity $\mathrm{d} s \otimes \mathrm{e}^{-u_{T-s}(\infty)r}\Pi(\mathrm{d} r) \otimes \mathrm{d} \nu$.
- ▶ \mathbb{N}_T^1 is a Poisson process on $[0, \infty)^2 \times \mathbb{N}_0$ with intensity $ds \otimes r e^{-u_{T-s}(\infty)r} \Pi(dr) \otimes \sharp(dj)$,
- ▶ $\mathbb{N}_T^2(ds, dr, dk, dj)$ is a Poisson process on $[0, \infty)^2 \times \mathbb{N}_0 \times \mathbb{N}$ with intensity

$$\left\{\frac{u_{T-s}(\infty)\psi'(u_{T-s}(\infty))-\psi(u_{T-s}(\infty))}{u_{T-s}(\infty)}\right\}\mathrm{d} s\otimes \eta_k^{T-s}(\mathrm{d} r)\otimes p_k^{T-s}\sharp(\mathrm{d} k)\otimes\sharp(\mathrm{d} j),$$

where, for $k \ge 2$,

$$\eta_k^{T-s}(\mathrm{d}r) = \frac{\beta u_{T-s}^2(\infty) \mathbf{1}_{\{k=2\}} \delta_0(\mathrm{d}r) + (u_{T-s}(\infty)r)^k \,\mathrm{e}^{-u_{T-s}(\infty)r} \Pi(\mathrm{d}r)/k!}{p_k^{T-s} (u_{T-s}(\infty)\psi'(u_{T-s}(\infty)) - \psi(u_{T-s}(\infty)))}, \qquad r \ge 0,$$

▶ W(ds, du) is the white noise process on $(0, \infty)^2$ based on the Lebesgue measure $ds \otimes du$.

CONDITIONING ON SURVIVAL.

The law of $(\Lambda_t^T, 0 \le t < T)$ conditional on $(\mathcal{F}^{\Lambda_t^T} \cap \{Z_0^T \ge 1\}, 0 \le t < T)$ is that of the law of the ψ -CSBP, *X*, conditioned to survive until time T.

This law is can be obtained by the following change of measure for $t \ge 0, x > 0$

$$\frac{\mathrm{d}\widetilde{\mathbb{P}}_x^T}{\mathrm{d}\mathbb{P}_x}\bigg|_{\mathcal{F}_t} = \frac{1 - \mathrm{e}^{-X_t u_{T-t}(\infty)}}{1 - \mathrm{e}^{-x u_T(\infty)}}.$$

• We have for
$$k \ge 1$$

$$\mathbf{P}_{x}^{T}[Z_{0}=k|Z_{0}\geq 1]=\frac{(u_{T}(\infty)x)^{k}}{k!}\frac{\mathrm{e}^{-u_{T}(\infty)x}}{1-\mathrm{e}^{-u_{T}(\infty)x}}.$$

▶ If n_T denotes the conditional probability $n(\cdot | \sup_{s \ge 0} \epsilon_s \ge T)$, then the first branch time γ_T of the individual corresponding to the excursion ϵ is given by

$$n_T(\gamma_T > t) = \frac{\psi(u_T(\infty))}{u_T(\infty)} \frac{u_{T-t}(\infty)}{\psi(u_{T-t}(\infty))},$$

for $t \in [0, T)$.

CONDITIONING ON SURVIVAL.

The law of $(\Lambda_t^T, 0 \le t < T)$ conditional on $(\mathcal{F}^{\Lambda_t^T} \cap \{Z_0^T \ge 1\}, 0 \le t < T)$ is that of the law of the ψ -CSBP, *X*, conditioned to survive until time T.

Take $T \to \infty$.

This law is can be obtained by the following change of measure for $t \ge 0, x > 0$

$$\frac{\mathrm{d}\widetilde{\mathbb{P}}_{x}^{T}}{\mathrm{d}\mathbb{P}_{x}}\bigg|_{\mathcal{F}_{t}} = \frac{1 - \mathrm{e}^{-X_{t}u_{T-t}(\infty)}}{1 - \mathrm{e}^{-xu_{T}(\infty)}} \longrightarrow \mathrm{e}^{-\alpha t}\frac{X_{t}}{x}$$

• We have for $k \ge 1$

$$\mathbf{P}_{x}^{T}[Z_{0}=k|Z_{0}\geq 1]=\frac{(u_{T}(\infty)x)^{k}}{k!}\frac{\mathrm{e}^{-u_{T}(\infty)x}}{1-\mathrm{e}^{-u_{T}(\infty)x}}\longrightarrow 0, \text{ unless } k=1.$$

▶ If n_T denotes the conditional probability $n(\cdot | \sup_{s \ge 0} \epsilon_s \ge T)$, then the first branch time γ_T of the individual corresponding to the excursion ϵ is given by

$$n_T(\gamma_T > t) = \frac{\psi(u_T(\infty))}{u_T(\infty)} \frac{u_{T-t}(\infty)}{\psi(u_{T-t}(\infty))} \longrightarrow 1,$$

for $t \in [0, T)$.

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Note that the convergence is in a weak sense.

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Theorem

Suppose that ψ is a critical or subcritical branching mechanism such that Grey's condition holds. Suppose, moreover, that $((\Lambda_t^T, Z_t^T), 0 \le t < T)$ is a weak solution to (3) and that Z_0^T is an independent random variable which is Poisson distributed with intensity $u_T(\infty)\Lambda_0^T$. Then, conditional on the event $Z_0^T > 0$, in the sense of weak convergence with respect to the Skorokhod topology on $\mathbb{D}([0,\infty), \mathbb{R}^2)$, for all t > 0,

$$((\Lambda_s^T, Z_s^T), 0 \le s \le t) \to ((X_s^{\uparrow}, 1), 0 \le s \le t),$$

where X^{\uparrow} is a weak solution to

$$\begin{aligned} X_t &= x + \alpha \int_0^t X_{s-} ds + \sqrt{2\beta} \int_0^t \int_0^{X_{s-}} W(ds, du) + \int_0^t \int_0^\infty \int_0^{X_{s-}} r \tilde{N}(ds, dr, du) \\ &+ \int_0^t \int_0^\infty r N^*(ds, dr) + 2\beta t, \qquad t \ge 0. \end{aligned}$$

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(DRIVING SOURCES OF RANDOMNESS III.)

In the previous theorem

- ▶ W(ds, du) is a white noise process on $(0, \infty)^2$ based on the Lebesgue measure $ds \otimes du$,
- ▶ $N(ds, dr, d\nu)$ is a Poisson point process on $(0, \infty)^3$ with intensity $ds \otimes \Pi(dr) \otimes d\nu$, and $\tilde{N}(ds, dr, d\nu)$ is the compensated measure of $N(ds, dr, d\nu)$,
- ▶ N^* is a Poisson random measure on $[0, \infty) \times (0, \infty)$ with intensity measure $ds \otimes r\Pi(dr)$.

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