# Meromorphic Lévy Processes and a new Wiener-Hopf Monte-Carlo simulation method

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- Other motivations from queuing theory, population models etc.
- One is fundamentally interested in the joint distribution

$$P(X_t \in dx, \, \overline{X}_t \in dy)$$

for any Lévy process (X, P).

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$$\Psi(\theta) := -\frac{1}{t} \log E(\mathrm{e}^{\mathrm{i}\theta X_t})$$
  
=  $a\mathrm{i}\theta + \frac{1}{2}\sigma^2\theta^2 + \int_{\mathbb{R}} (1 - \mathrm{e}^{\mathrm{i}\theta x} + \mathrm{i}\theta x \mathbf{1}_{\{|x| \le 1\}}) \Pi(dx)$ 

where  $a \in \mathbb{R}$ ,  $\sigma \in \mathbb{R}$  and  $\Pi$  is a measure concentrated on  $\mathbb{R} \setminus \{0\}$  satisfying  $\int_{\mathbb{R}} (1 \wedge x^2) \Pi(dx) < \infty...$ 

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$$\begin{split} \Psi(\theta) &:= -\frac{1}{t} \log E(\mathrm{e}^{\mathrm{i}\theta X_t}) \\ &= a\mathrm{i}\theta + \frac{1}{2}\sigma^2\theta^2 + \int_{\mathbb{R}} (1 - \mathrm{e}^{\mathrm{i}\theta x} + \mathrm{i}\theta x \mathbf{1}_{\{|x| \le 1\}}) \Pi(dx) \end{split}$$

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- ...in which case there are fast-Fourier methods for inverting  $\exp\{-\Psi(\theta)t\}$  to give  $P(X_t \in dx)$ .
- For the case of  $P(\overline{X}_t \in dx)$ , recent methods have concentrated on Fourier inversion of the, so called, Wiener-Hopf factors.

Recall that it turns out that one may always uniquely decompose

$$\frac{q}{q+\Psi(\theta)} = E(\mathrm{e}^{\mathrm{i}\theta \overline{X}_{\mathbf{e}_q}}) \times E(\mathrm{e}^{\mathrm{i}\theta \underline{X}_{\mathbf{e}_q}})$$

where  $e_q$  is an independent and exponentially distributed random variable with rate q > 0 and  $\underline{X}_t = \inf_{s \le t} X_s$ .

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• If one is in possession of close analytical expressions for these factors, Fourier inversion, first in  $\theta$  and then in q would be an option for accessing the law of  $\overline{X}_t$  and  $\underline{X}_t$ . However one is rarely in possession of the factors (even after 60 years of research into this topic), and even then there is the issue of the double Fourier inversion.

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- There are no convenient formulae which contain both  $X_t$  and  $\overline{X}_t$  which could be Fourier inverted.

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- Note that

$$\mathbf{g}(n,q) := \sum_{i=1}^{n} \frac{1}{q} \mathbf{e}^{(i)}$$

is a Gamma (Erlang) distribution with parameters n and q and by the strong law of Large numbers, for t > 0,

$$\mathbf{g}(n, n/t) = \sum_{i=1}^{n} \frac{t}{n} \mathbf{e}^{(i)} \to t$$

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• Hence for a suitably large n, we have in distribution

$$(X_{\mathbf{g}(n,n/t)}, \overline{X}_{\mathbf{g}(n,n/t)}) \simeq (X_t, \overline{X}_t).$$

Indeed since t is not a jump time with probability 1, we have that  $(X_{\mathbf{g}(n,n/t)}, \overline{X}_{\mathbf{g}(n,n/t)}) \rightarrow (X_t, \overline{X}_t)$  almost surely.

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A reformulation of the Wiener-Hopf factorization states that

$$X_{\mathbf{e}_q} \stackrel{d}{=} S_q + I_q$$

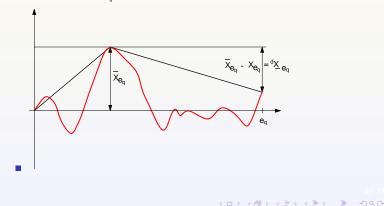
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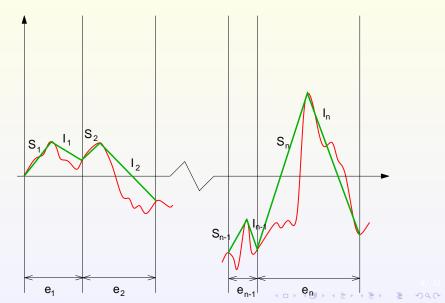


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• Taking advantage of the above, the fact that X has stationary and independent increments and the fact that, as a time, g(n, n/t) can be seen as n subsequent independent exponential time periods we have the following:



• Theorem. For all 
$$n \in \{1, 2, \dots\}$$
 and  $q > 0$ ,

$$(X_{\mathbf{g}(n,q)}, \overline{X}_{\mathbf{g}(n,q)}) \stackrel{d}{=} (V(n,q), J(n,q))$$

where

$$V(n,q) = \sum_{j=1}^{n} \{S_q^{(j)} + I_q^{(j)}\} \text{ and } J(n,q) := \bigvee_{i=0}^{n-1} \left(\sum_{j=1}^{i} \{S_q^{(j)} + I_q^{(j)}\} + S_q^{(i+1)}\right)$$

Here,  $S_q^{(0)} = I_q^{(0)} = 0$ ,  $\{S_q^{(j)} : j \ge 1\}$  are an i.i.d. sequence of random variables with common distribution equal to that of  $\overline{X}_{\mathbf{e}_q}$  and  $\{I_q^{(j)} : j \ge 1\}$  are another i.i.d. sequence of random variable with common distribution equal to that of  $\underline{X}_{\mathbf{e}_q}$ .

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- Moreover, we have the following obvious:
- **Corollary.** We have as  $n \uparrow \infty$

$$(V(n, n/t), J(n, n/t)) \rightarrow (X_t, \overline{X}_t)$$

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where the convergence is understood in the distributional sense.

The previous results suggest that to simulate, for example,  $\mathbb{E}(g(X_t, \overline{X}_t))$  one should follow the following algorithm:

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Then approximate

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- This numerical procedure has disposed of one (numerical) Fourier inverse computation.
- This still leaves the problem of simulating from the unknown distribution  $\overline{X}_{\mathbf{e}_{n/t}}$  and  $\underline{X}_{\mathbf{e}_{n/t}}$  i.e. we are still one (numerical) Fourier transform away from  $(X_t, \overline{X}_t)$

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# Meromorphic Lévy processes (definition)

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A Lévy process is said to belong to the Meromorphic class (*M*-class), if and only if the Lévy measure II(dx) has a density with respect to the Lebesgue measure, given by

$$\pi(x) = \mathbb{I}_{\{x>0\}} \sum_{n\geq 1} a_n \rho_n e^{-\rho_n x} + \mathbb{I}_{\{x<0\}} \sum_{n\geq 1} \hat{a}_n \hat{\rho}_n e^{\hat{\rho}_n x}, \tag{1}$$

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where all the coefficients  $a_n$ ,  $\hat{a}_n$ ,  $\rho_n$ ,  $\hat{\rho}_n$  are positive, the sequences  $\{\rho_n\}_{n\geq 1}$  and  $\{\hat{\rho}_n\}_{n\geq 1}$  are strictly increasing, and  $\rho_n \to +\infty$  and  $\hat{\rho}_n \to +\infty$  as  $n \to +\infty$ .

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- We allow the case of a finite number summands (on either or both sides of the origin) with obvious modifications to the above.
- $\blacksquare$  To ensure that  $\int_{\mathbb{R}}x^2\pi(x)\mathrm{d}x$  converges we need to impose the additional constraint that

$$\sum_{n \ge 1} a_n \rho_n^{-2} + \sum_{n \ge 1} \hat{a}_n \hat{\rho}_n^{-2} < \infty$$

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#### Meromorphic Lévy processes (equivalent definition)

- (i) The characteristic exponent  $\Psi(z)$  is a meromorphic function which has poles at points  $\{-i\rho_n, i\hat{\rho}_n\}_{n\geq 1}$ , where  $\rho_n$  and  $\hat{\rho}_n$  are positive real numbers.
- (ii) For q ≥ 0 function q + Ψ(z) has roots at points {-iζ<sub>n</sub>, iζ̂<sub>n</sub>}<sub>n≥1</sub> where ζ<sub>n</sub> and ζ̂<sub>n</sub> are nonnegative real numbers (strictly positive if q > 0). We will write ζ<sub>n</sub>(q), ζ̂<sub>n</sub>(q) if we need to stress the dependence on q.
- (iii) The roots and poles of  $q + \Psi(iz)$  satisfy the following interlacing condition

$$\dots - \rho_2 < -\zeta_2 < -\rho_1 < -\zeta_1 < 0 < \hat{\zeta}_1 < \hat{\rho}_1 < \hat{\zeta}_2 < \hat{\rho}_2 < \dots$$

(iv) The Wiener-Hopf factors are expressed as convergent infinite products,

$$\mathbb{E}\left[\mathrm{e}^{-z\overline{X}_{\mathbf{e}_{q}}}\right] = \prod_{n\geq 1} \frac{1+\frac{z}{\rho_{n}}}{1+\frac{z}{\zeta_{n}}}$$
$$\mathbb{E}\left[\mathrm{e}^{z\underline{X}_{\mathbf{e}_{q}}}\right] = \prod_{n\geq 1} \frac{1+\frac{z}{\rho_{n}}}{1+\frac{z}{\zeta_{n}}}.$$

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The density of the Lévy measure is

$$\pi(x) = \mathbf{1}_{\{x>0\}} \sum_{i=1}^{N} a_i \rho_i \mathrm{e}^{-\rho_i x} + \mathbf{1}_{\{x<0\}} \sum_{i=1}^{\hat{N}} \hat{a}_i \hat{\rho}_i \mathrm{e}^{\hat{\rho}_i x},$$

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Including Gaussian and linear drift, one can verify that the characteristic exponent is a rational function and that hyper-exponential Lévy processes have finite activity jumps and paths of bounded variation unless σ > 0.

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- Including Gaussian and linear drift, one can verify that the characteristic exponent is a rational function and that hyper-exponential Lévy processes have finite activity jumps and paths of bounded variation unless σ > 0.
- Note that this class has been looked at by many other authors in the past and historically is starts life as the Kou process.

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## **Example:** Kuznetsov's $\beta$ -family

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## **Example:** Kuznetsov's $\beta$ -family

• The characteristic exponent  $(\Psi(\theta) = -\log \mathbb{E}(e^{i\theta X_1}), \theta \in \mathbb{R})$  is given by

$$\Psi(\theta) = iaz + \frac{1}{2}\sigma^2 z^2 + \frac{c_1}{\beta_1} \left\{ \mathsf{B}(\alpha_1, 1 - \lambda_1) - \mathsf{B}(\alpha_1 - \frac{\mathrm{i}\theta}{\beta_1}, 1 - \lambda_1) \right\} \\ + \frac{c_2}{\beta_2} \left\{ \mathsf{B}(\alpha_2, 1 - \lambda_2) - \mathsf{B}(\alpha_2 + \frac{\mathrm{i}\theta}{\beta_2}, 1 - \lambda_2) \right\}$$

where  $B(x, y) = \Gamma(x)\Gamma(y)/\Gamma(x + y)$  is the Beta function, with parameter range  $a \in \mathbb{R}, \sigma, c_i, \alpha_i, \beta_i > 0$  and  $\lambda_1, \lambda_2 \in (0, 3) \setminus \{1, 2\}$ .

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where  $\mathsf{B}(x,y) = \Gamma(x)\Gamma(y)/\Gamma(x+y)$  is the Beta function, with parameter range  $a \in \mathbb{R}, \sigma, c_i, \alpha_i, \beta_i > 0$  and  $\lambda_1, \lambda_2 \in (0,3) \setminus \{1,2\}$ .

■ The corresponding Lévy measure Π has density

$$\pi(x) = c_1 \frac{\mathrm{e}^{-\alpha_1 \beta_1 x}}{(1 - \mathrm{e}^{-\beta_1 x})^{\lambda_1}} \mathbf{1}_{\{x > 0\}} + c_2 \frac{\mathrm{e}^{\alpha_2 \beta_2 x}}{(1 - \mathrm{e}^{\beta_2 x})^{\lambda_2}} \mathbf{1}_{\{x < 0\}}.$$

The  $\beta$ -class of Lévy processes includes another recently introduced family of Lévy processes known as Lamperti-stable processes.

## **Example: Hypergeometric Lévy processes**

# Example: Hypergeometric Lévy processes

• The characteristic exponent (  $\Psi(\theta)=\mathbb{E}(e^{i\theta X_1}), \theta\in\mathbb{R})$  is given by

$$\Psi(\theta) = \frac{\Gamma(1 - \beta + \gamma - i\theta)}{\Gamma(1 - \beta + i\theta)} \frac{\Gamma(\hat{\beta} + \hat{\gamma} + i\theta)}{\Gamma(\hat{\beta} + i\theta)}$$

where  $(\beta,\gamma,\hat{\beta},\hat{\gamma})$  belong to the admissible range

$$\{\beta \le 1, \gamma \in (0,1), \hat{\beta} \ge 0, \hat{\gamma} \in (0,1)\}.$$

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where  $(\beta,\gamma,\hat{\beta},\hat{\gamma})$  belong to the admissible range

$$\{\beta \le 1, \gamma \in (0,1), \hat{\beta} \ge 0, \hat{\gamma} \in (0,1)\}.$$

The Lévy density is given by

$$\pi(x) = \begin{array}{c} -\frac{\Gamma(\eta)}{\Gamma(\eta-\hat{\gamma})\Gamma(-\gamma)} \mathrm{e}^{-(1-\beta+\gamma)x} {}_2F_1(1+\gamma,\eta;\eta-\hat{\gamma};\mathrm{e}^{-x}) & \text{if } x > 0\\ -\frac{\Gamma(\eta)}{\Gamma(\eta-\gamma)\Gamma(-\hat{\gamma})} \mathrm{e}^{(\hat{\beta}+\hat{\gamma})x} {}_2F_1(1+\hat{\gamma},\eta;\eta-\gamma;\mathrm{e}^{-x}) & \text{if } x < 0 \end{array}$$

where  $\eta = 1 - \beta + \gamma + \hat{\beta} + \hat{\gamma}$ .

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# **Distribution of extrema**

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For  $x \ge 0$ 

$$\mathbb{P}(\overline{X}_{\mathbf{e}_q} \in \mathsf{d}x) = \bar{\mathbf{a}}(\rho, \zeta)^T \times \bar{\mathbf{v}}(\zeta, x) \mathsf{d}x \\ \mathbb{P}(-\underline{X}_{\mathbf{e}_q} \in \mathsf{d}x) = \bar{\mathbf{a}}(\hat{\rho}, \hat{\zeta})^T \times \bar{\mathbf{v}}(\hat{\zeta}, x) \mathsf{d}x.$$

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#### **Distribution of extrema**

• For  $x \ge 0$  $\mathbb{P}(\overline{X}_{e_q} \in dx) = \bar{\mathbf{a}}(\rho, \zeta)^T \times \bar{\mathbf{v}}(\zeta, x) dx$   $\mathbb{P}(-\underline{X}_{e_q} \in dx) = \bar{\mathbf{a}}(\hat{\rho}, \hat{\zeta})^T \times \bar{\mathbf{v}}(\hat{\zeta}, x) dx.$ • Here  $\bar{\mathbf{a}}(\rho, \zeta) = [\mathbf{a}_0(\rho, \zeta), \mathbf{a}_1(\rho, \zeta), \mathbf{a}_2(\rho, \zeta), \ldots]^T$  such that  $\mathbf{a}_0(\rho, \zeta) = \lim_{n \to +\infty} \prod_{k=1}^n \frac{\zeta_k}{\rho_k}, \quad \mathbf{a}_n(\rho, \zeta) = \left(1 - \frac{\zeta_n}{\rho_n}\right) \prod_{\substack{k \ge 1 \\ k \ne n}} \frac{1 - \frac{\zeta_n}{\rho_k}}{1 - \frac{\zeta_n}{\zeta_k}}$ 

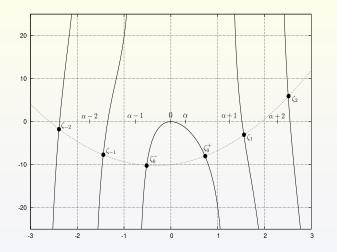
$$\bar{\mathbf{v}}(\zeta, x) = \left[\delta_0(x), \zeta_1 e^{-\zeta_1 x}, \zeta_2 e^{-\zeta_2 x}, \dots\right]^T$$

where  $\delta_0(x)$  is the Dirac delta function at x = 0. A similar expression holds for  $\bar{\mathbf{a}}(\rho, \zeta)$ .

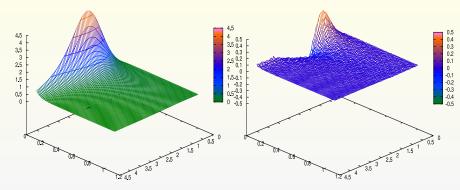
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# **Computing roots**



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**Figure:** Computing the joint density of  $(\overline{X}_1, \overline{X}_1 - X_1)$  for parameter Set 1. Here  $\overline{X}_1 \in [0, 1]$  and  $\overline{X}_1 - X_1 \in [0, 4]$ . Fourier method benchmark (left), N = 20, WH-MC error (right).

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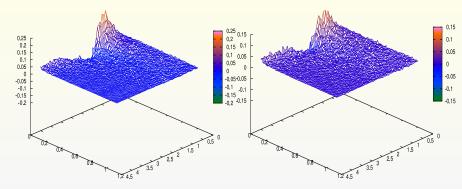
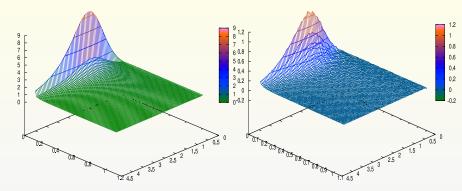


Figure: Computing the joint density of  $(\overline{X}_1, \overline{X}_1 - X_1)$  for parameter Set 1. Here  $\overline{X}_1 \in [0, 1]$  and  $\overline{X}_1 - X_1 \in [0, 4]$ . N = 50, WH-MC error(left), N = 100, WH-MC error (right).

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**Figure:** Computing the joint density of  $(\overline{X}_1, \overline{X}_1 - X_1)$  for parameter Set 1. Here  $\overline{X}_1 \in [0, 1]$  and  $\overline{X}_1 - X_1 \in [0, 4]$ . N = 100, MC simulation (left), N = 100, MC error (right).

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Total computation time for WH-MC is at most half the computation time for Fourier inversion of  $\exp -\Psi(z)$  followed by a random walk simulation.

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- The overwhelming majority of the WH-MC method is the simulation, computing the roots takes 1% of the time. Roots can be stored once they have been computed.
- Considerably more accurate for the same number of steps in each cycle.
- Does not artificially build in an atom at zero in the numerical distribution of  $\overline{X}_t$ .

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## More identities: One sided exit problem

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• Define a matrix  $\mathbf{A} = \{a_{i,j}\}_{i,j \ge 0}$  as

$$a_{i,j} = \begin{cases} 0 & \text{if } i = 0, \ j \ge 0\\ a_i(\rho, \zeta) b_0(\zeta, \rho) & \text{if } i \ge 1, \ j = 0\\ \frac{a_i(\rho, \zeta) b_j(\zeta, \rho)}{\rho_j - \zeta_i} & \text{if } i \ge 1, \ j \ge 1 \end{cases}$$
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Then for c > 0 and  $y \ge 0$  we have

$$\mathbb{E}\left[e^{-q\tau_c^+}\mathbb{I}\left(X_{\tau_c^+} - c \in \mathsf{d}y\right)\right] = \bar{\mathsf{v}}(\zeta, c)^T \times \mathbf{A} \times \bar{\mathsf{v}}(\rho, y)\mathsf{d}y.$$
(3)

#### More identities: Two-sided exit problem

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• Let a > 0 and define a matrix  $\mathbf{B} = \mathbf{B}(\hat{\rho}, \zeta, a) = \{b_{i,j}\}_{i,j \ge 0}$  with

$$b_{i,j} = \begin{cases} \zeta_j e^{-a\zeta_j} & \text{if } i = 0, \ j \ge 1\\ 0 & \text{if } i \ge 0, \ j = 0\\ \frac{\hat{\rho}_i \zeta_j}{\hat{\rho}_i + \zeta_j} e^{-a\zeta_j} & \text{if } i \ge 1, \ j \ge 1 \end{cases}$$

and similarly  $\hat{\mathbf{B}} = \mathbf{B}(\rho, \hat{\zeta}, a)$ . There exist matrices  $\mathbf{C}_1$ ,  $\mathbf{C}_2$  such that for  $x \in (0, a)$  we have

$$\begin{split} \mathbb{E}_{x} \left[ \mathrm{e}^{-q\tau_{a}^{+}} \mathbb{I} \left( X_{\tau_{a}^{+}} \in \mathsf{d}y \; ; \; \tau_{a}^{+} < \tau_{0}^{-} \right) \right] \\ &= \left[ \bar{\mathsf{v}}(\zeta, a - x)^{T} \times \mathbf{C}_{1} + \bar{\mathsf{v}}(\hat{\zeta}, x)^{T} \times \mathbf{C}_{2} \right] \times \bar{\mathsf{v}}(\rho, y - a) \mathsf{d}y \end{split}$$

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These matrices satisfy the following system of linear equations

$$\begin{cases} \mathbf{C}_1 &= \mathbf{A} - \hat{\mathbf{C}}_2 \mathbf{B} \mathbf{A} \\ \hat{\mathbf{C}}_2 &= -\mathbf{C}_1 \hat{\mathbf{B}} \hat{\mathbf{A}} \end{cases} \qquad \begin{cases} \hat{\mathbf{C}}_1 &= \hat{\mathbf{A}} - \mathbf{C}_2 \hat{\mathbf{B}} \hat{\mathbf{A}} \\ \mathbf{C}_2 &= -\hat{\mathbf{C}}_1 \mathbf{B} \mathbf{A} \end{cases}$$

This system of linear equations can be solved iteratively with exponential convergence.

• For a > 0,  $y \le a$  we define  $R^{(q)}(a, dy) := \int_0^\infty e^{-qt} \mathbb{P}(X_t \in dy; t < \tau_a^+) dt$ 

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- For a > 0,  $y \le a$  we define  $R^{(q)}(a, dy) := \int_0^\infty e^{-qt} \mathbb{P}(X_t \in dy; t < \tau_a^+) dt$
- Define a matrix  $\mathbf{D} = \{d_{i,j}\}_{i,j\geq 0}$  as follows

$$d_{i,j} = \begin{cases} 0 & \text{if } i = 0 \text{ or } j = 0\\ \mathsf{a}_i(\rho,\zeta) \frac{\zeta_i \hat{\zeta}_j}{\zeta_i + \hat{\zeta}_j} \mathsf{a}_j(\hat{\rho},\hat{\zeta}) & \text{if } i \ge 1, \ j \ge 1 \end{cases}$$

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- For a > 0,  $y \le a$  we define  $R^{(q)}(a, \mathrm{d}y) := \int_0^\infty \mathrm{e}^{-qt} \mathbb{P}(X_t \in \mathrm{d}y; t < \tau_a^+) \mathrm{d}t$
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• Then if  $y \leq a$  we have

$$\begin{split} qR^{(q)}(a,\mathrm{d}y) &= \left[\bar{\mathsf{v}}(\zeta,0\vee y)\times\mathbf{D}\times\bar{\mathsf{v}}(\hat{\zeta},0\vee(-y))\right)\\ &-\bar{\mathsf{v}}(\zeta,a)\times\mathbf{D}\times\bar{\mathsf{v}}(\hat{\zeta},a-y)\big]\mathrm{d}y. \end{split}$$

# **Example of numerics**

## **Example of numerics**

Choose an example from Kuznetsov's β-class that has bounded variation jump component and concentrate on four cases: With/without Gaussian component, drift to ±∞.

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## **Example of numerics**

- Choose an example from Kuznetsov's β-class that has bounded variation jump component and concentrate on four cases: With/without Gaussian component, drift to ±∞.
- For the above four cases, consider the following densities.
  - (i) density of the overshoot if the exit happens at the upper boundary

$$f_1(x,y) = \frac{\mathsf{d}}{\mathsf{d}y} \mathbb{E}_x \left[ \mathrm{e}^{-q\tau_1^+} \mathbb{I}\left( X_{\tau_1^+} \leq y \ ; \ \tau_1^+ < \tau_0^- \right) \right]$$

(ii) probability of exiting from the interval  $\left[0,1\right]$  at the upper boundary

$$f_2(x) = \mathbb{E}_x \left[ e^{-q\tau_1^+} \mathbb{I} \left( \tau_1^+ < \tau_0^- \right) \right]$$

(iii) probability of exiting the interval  $\left[0,1\right]$  by creeping across the upper boundary

$$f_3(x) = \mathbb{E}_x \left[ e^{-q\tau_1^+} \mathbb{I} \left( X_{\tau_1^+} = 1 \; ; \; \tau_1^+ < \tau_0^- \right) \right]$$

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Double sided exit:  $\sigma > 0$  and positive drift

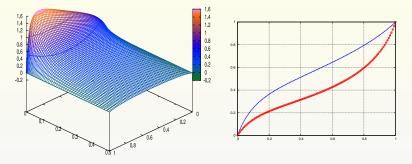


Figure: Unbounded variation case ( $\sigma = 0.5$ ): computing the density of the overshoot  $f_1(x, y)$  ( $x \in (0, 1)$ ,  $y \in (0, 0.5)$ ), probability of first exit  $f_2(x)$  and probability of creeping  $f_3(x)$  for parameter Set 1, positive drift  $\mu = 1$ 

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#### Double sided exit: $\sigma > 0$ and negative drift

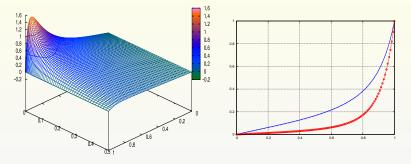
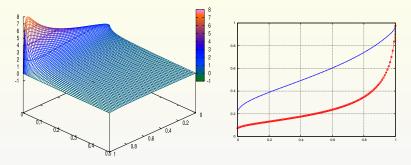


Figure: Unbounded variation case ( $\sigma = 0.5$ ): computing the density of the overshoot  $f_1(x, y)$  ( $x \in (0, 1)$ ,  $y \in (0, 0.5)$ ), probability of first exit  $f_2(x)$  and probability of creeping  $f_3(x)$  for parameter Set 2, negative drift  $\mu = -1$ .

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#### Double sided exit: bounded variation and positive drift

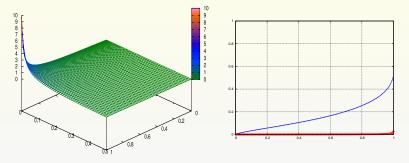


**Figure:**  $\sigma = 0$ , positive drift: computing the density of the overshoot  $f_1(x, y)$   $(x \in (0, 1), y \in (0, 0.5))$ , probability of first exit  $f_2(x)$  and probability of creeping  $f_3(x)$ .

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#### Double sided exit: bounded variation and negative drift



**Figure:**  $\sigma = 0$ , negative drift: computing the density of the overshoot  $f_1(x, y)$   $(x \in (0, 1), y \in (0, 0.5))$ , probability of first exit  $f_2(x)$  and probability of creeping  $f_3(x)$ .

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## Simulating processes with heavy tails

#### Simulating processes with heavy tails

 A little thought shows that a huge class of Lévy processes can be written as the independent sum of a β-process plus and independent compound Poisson process. Say,

$$Y_t = X_t + \sum_{i=1}^{N_t} \xi_i$$

where  $\{N_t:t\geq 0\}$  is a Poisson process of rate  $\gamma$  and  $\{\xi_i:i\geq 1\}$  and i.i.d. sequence.

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where  $\{N_t:t\geq 0\}$  is a Poisson process of rate  $\gamma$  and  $\{\xi_i:i\geq 1\}$  and i.i.d. sequence.

 $\blacksquare$  Define iteratively for  $n\geq 1$ 

$$V(n,\lambda) = V(n-1,\lambda) + S^{(n)}_{\lambda+\gamma} + I^{(n)}_{\lambda+\gamma} + \xi_n(1-\beta_n)$$
  
$$J(n,\lambda) = \max\left(V(n,\lambda), J(n-1,\lambda), V(n-1,\lambda) + S^{(n)}_{\lambda+\gamma}\right)$$

where  $V(0, \lambda) = J(0, \lambda) = 0$  and  $\{\beta_n : n \ge 1\}$  are an i.i.d. sequence of Bernoulli random variables such that  $\mathbb{P}(\beta_n = 1) = \lambda/(\gamma + \lambda)$ . Then

$$(Y_{\mathbf{g}(n,\lambda)}, \overline{Y}_{\mathbf{g}(n,\lambda)}) \stackrel{d}{=} (V(T_n,\lambda), J(T_n,\lambda))$$
  
where  $T_n = \min\{j \ge 1 : \sum_{i=1}^{j} \beta_i = n\}.$ 

# Approximate simulation of the law of $(X_t, \overline{X}_t, \underline{X}_t)$

Define iteratively for  $n \ge 1$ 

$$\begin{split} V(n,\lambda) &= V(n-1,\lambda) + S_{\lambda}^{(n)} + I_{\lambda}^{(n)} \\ J(n,\lambda) &= \max\left(J(n-1,\lambda), V(n-1,\lambda) + S_{\lambda}^{(n)}\right) \\ K(n,\lambda) &= \min\left(K(n-1,\lambda), V(n,\lambda)\right) \\ \tilde{J}(n,\lambda) &= \max\left(\tilde{J}(n-1,\lambda), V(n,\lambda)\right) \\ \tilde{K}(n,\lambda) &= \min\left(\tilde{K}(n-1,\lambda), V(n-1,\lambda) + I_{\lambda}^{(n)}\right), \end{split}$$

where  $V(0, \lambda) = J(0, \lambda) = K(0, \lambda) = \tilde{J}(0, \lambda) = \tilde{K}(0, \lambda) = 0$ . Then for any bounded function  $f(x, y, z) : \mathbb{R}^3 \to \mathbb{R}$  which is increasing in z-variable we have

$$\mathbb{E}[f(V(n,\lambda), J(n,\lambda), K(n,\lambda))] \geq \mathbb{E}[f(X_{\mathbf{g}(n,\lambda)}, \overline{X}_{\mathbf{g}(n,\lambda)}, \underline{X}_{\mathbf{g}(n,\lambda)})] \\ \mathbb{E}[f(V(n,\lambda), \tilde{K}(n,\lambda), \tilde{J}(n,\lambda))] \leq \mathbb{E}[f(X_{\mathbf{g}(n,\lambda)}, \underline{X}_{\mathbf{g}(n,\lambda)}, \overline{X}_{\mathbf{g}(n,\lambda)})].$$

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