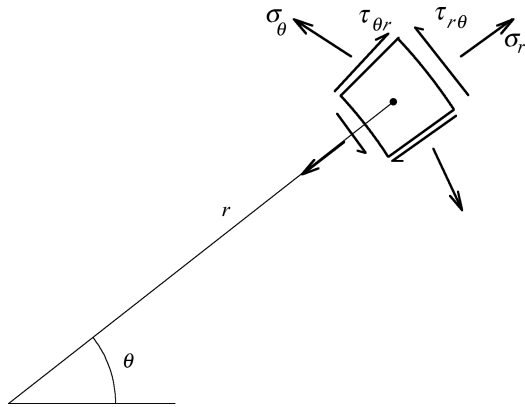


Airy stress function in Polar coordinates



The diagram shows a plan view of an element of a shell in cylindrical polar coordinates, r, θ, z . The horizontal components of membrane stress per unit horizontal length are $\sigma_r, \sigma_\theta, \tau_{r\theta}$ and $\tau_{\theta r}$. If there is no horizontal load, the equilibrium equations in plan are

$$\frac{\partial}{\partial r}(r\sigma_r) - \sigma_\theta + \frac{\partial \tau_{\theta r}}{\partial \theta} = 0, \quad \frac{\partial \sigma_\theta}{\partial \theta} + \frac{\partial}{\partial r}(r\tau_{r\theta}) + \tau_{\theta r} = 0 \quad \text{and} \quad \tau_{\theta r} = \tau_{r\theta}.$$

They come from considering element $\delta r \delta \theta$:

Resolving in the radial direction:

$$\frac{\partial}{\partial r}(r\delta\theta\sigma_r)\delta r - 2\delta r\sigma_\theta\frac{\delta\theta}{2} + \frac{\partial}{\partial \theta}(\delta r\tau_{\theta r})\delta\theta = 0.$$

Resolving in the tangential direction:

$$\frac{\partial}{\partial \theta}(\delta r\sigma_\theta)\delta\theta + \frac{\partial}{\partial r}(r\delta\theta\tau_{r\theta})\delta r + 2\delta r\tau_{\theta r}\frac{\delta\theta}{2} = 0.$$

Equilibrium of moments:

$$r\delta\theta\tau_{\theta r}\delta r = \delta r\tau_{r\theta}r\delta\theta.$$

The equilibrium equations in plan are automatically satisfied by $\sigma_r = \frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2}$,

$\sigma_\theta = \frac{\partial^2 \phi}{\partial r^2}$ and $\tau_{r\theta} = \tau_{\theta r} = -\frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \phi}{\partial \theta} \right)$ in which ϕ is an Airy stress function.

Proof:

$$\begin{aligned}
& \frac{\partial}{\partial r}(r\sigma_r) - \sigma_\theta + \frac{\partial \tau_{\theta r}}{\partial \theta} \\
&= \frac{\partial}{\partial r} \left(\frac{\partial \phi}{\partial r} + \frac{1}{r} \frac{\partial^2 \phi}{\partial \theta^2} \right) - \frac{\partial^2 \phi}{\partial r^2} - \frac{\partial^2}{\partial \theta \partial r} \left(\frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) \\
&= \frac{\partial^2 \phi}{\partial r^2} + \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial^2 \phi}{\partial \theta^2} \right) - \frac{\partial^2 \phi}{\partial r^2} - \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial^2 \phi}{\partial \theta^2} \right) = 0
\end{aligned}$$

$$\begin{aligned}
& \frac{\partial \sigma_\theta}{\partial \theta} + \frac{\partial}{\partial r}(r\tau_{r\theta}) + \tau_{\theta r} \\
&= \frac{\partial}{\partial \theta} \left(\frac{\partial^2 \phi}{\partial r^2} \right) - \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) \right) - \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) \\
&= \frac{\partial^3 \phi}{\partial \theta \partial r^2} - \frac{\partial}{\partial r} \left(\frac{\partial^2 \phi}{\partial \theta \partial r} - \frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) - \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) = 0
\end{aligned}$$

If z is the height of the shell, the vertical equilibrium equation is

$$w = \left(\frac{1}{r} \frac{\partial \phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} \right) \frac{\partial^2 z}{\partial r^2} - 2 \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial z}{\partial \theta} \right) + \left(\frac{1}{r} \frac{\partial z}{\partial r} + \frac{1}{r^2} \frac{\partial^2 z}{\partial \theta^2} \right) \frac{\partial^2 \phi}{\partial r^2}$$

in which w is downwards the load per unit plan area.

For radial symmetry all the derivatives with respect to θ disappear:

$$w = \frac{1}{r} \frac{d\phi}{dr} \frac{d^2 z}{dr^2} + \frac{1}{r} \frac{dz}{dr} \frac{d^2 \phi}{dr^2} = \frac{1}{r} \frac{d}{dr} \left(\frac{d\phi}{dr} \frac{dz}{dr} \right).$$